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Exotic spin states in frustrated antiferromagnet $MnSc_2S_4$

ESS-ILL User Meeting, Lund, 6 October, 2022



Main collaborators

Shang Gao (PhD 2017) and Guratinder Kaur (PhD 2021) (PSI Villigen, Switzerland)

Diego Rosales and Flavia Gomez Albarracin (La Plata, Argentina)

Vladimir Tsurkan (Augsburg, Germany; Chisinau Moldova)

Local contacts at Neutron Scattering Facilities:

T. Fennell, J. White, G. S. Tucker, B. Roessli, R. Sibille, LNS, SINQ, Villigen

J. Su, P. Cermak, A. Schneidewind, FRM2, Munich

E. Ressouche, S. Raymond, M. Boehm, P. Steffens, F. Bourdarot, ILL, Grenoble

Publications:

S. Gao et al. Nature Physics 13, 157 2017

S. Gao et al. Nature 583, 37 2020

D. Rosales et al. PhysRevB 105, 224402 2022

K. Guratinder et al. PhysRevB 105, 174422 2022



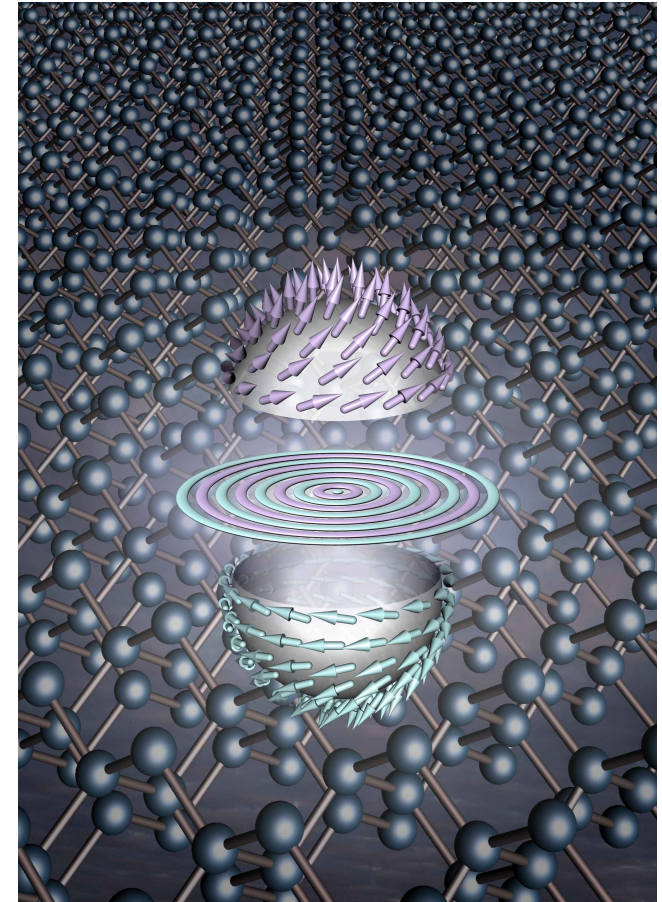
FONDS NATIONAL SUISSE
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FONDO NAZIONALE SVIZZERO
SWISS NATIONAL SCIENCE FOUNDATION

Spiral spin liquids

Multi-step ordering in MnSC_2S_4

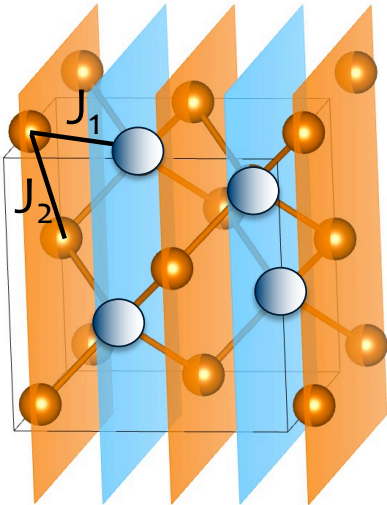
Exchange Hamiltonian of MnSC_2S_4

3k-structure and fractional AF skyrmion lattice



Diego Rosales, Private Collection

Spiral spin liquid on diamond lattice



$$H = J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \mathbf{S}_j + J_2 \sum_{\langle\langle ij \rangle\rangle} \mathbf{S}_i \mathbf{S}_j$$

J_2 - AF, introduces frustration (fcc)

J_1 - can increase frustration

two fcc sublattices

$T=0$: Ordering vectors lie on continuous surface in momentum space

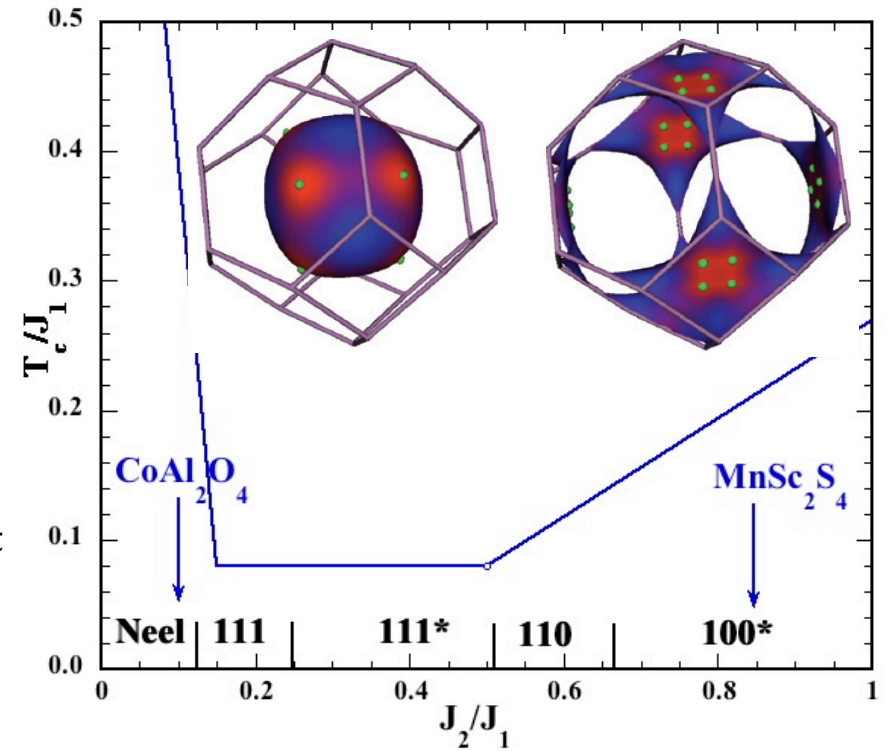
massive spiral degeneracy

spins can distort from GS1 to GS2 with no E cost

entropy stabilizes LRO at finite T by lifting degeneracy of the free energy along the spiral surface

order-by-disorder

D. Bergman et al. Nature Physics 3, 487(2007)



Spin spirals on bipartite lattices

T. A. Kaplan *Phys. Rev.* 116, 888(1959)

$$H = J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \mathbf{S}_j + J_2 \sum_{\langle\langle ij \rangle\rangle} \mathbf{S}_i \mathbf{S}_j$$

Lattices

square C. L. Henley *PRL* 62, 2056 1989

diamond S. Gao et al. *Nature Physics* 13, 157 2017 MnSc_2S_4

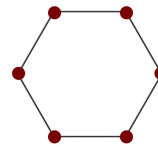
honeycomb A. Mulder et al. *PRB* 81, 214419(2010)
S. Gao et al. *PRL* 128, 227201(2022) FeCl_3

Correspondence between

Spin spiral surface of frustrated AF with

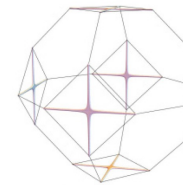
Fermi surface of noninteracting fermion systems

triangular lattice



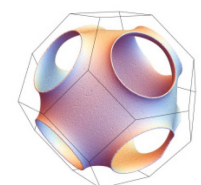
Dirac points

FCC lattice



nodal lines

diamond lattice

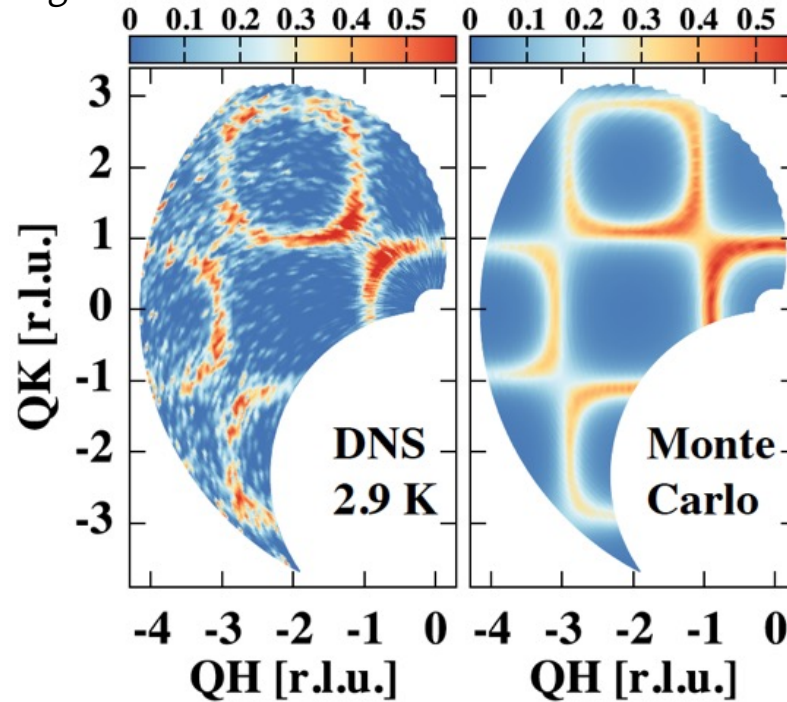


Fermi surface

J. Attig, S. Trebst *PRB* 96, 085145(2017)

AF spinel MnSc₂S₄ and spiral spin liquid

Neutron diffuse scattering



Classical MC calculation

8x10x10x10 atoms $S=5/2$

J_1 - J_2 Heisenberg Hamiltonian

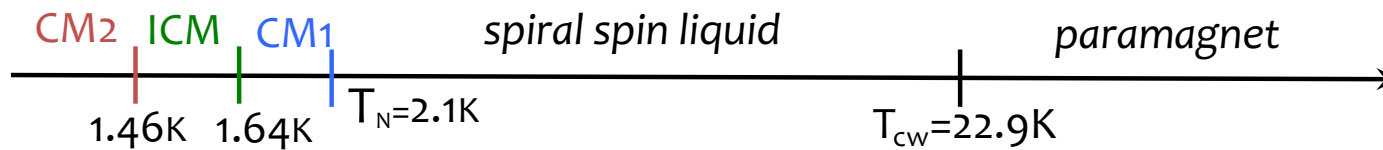
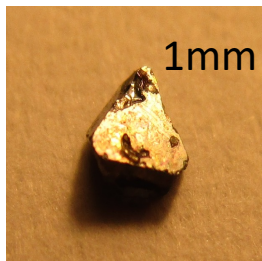
$$|J_2/J_1|=0.85$$

$$T/|J_1|=0.55$$

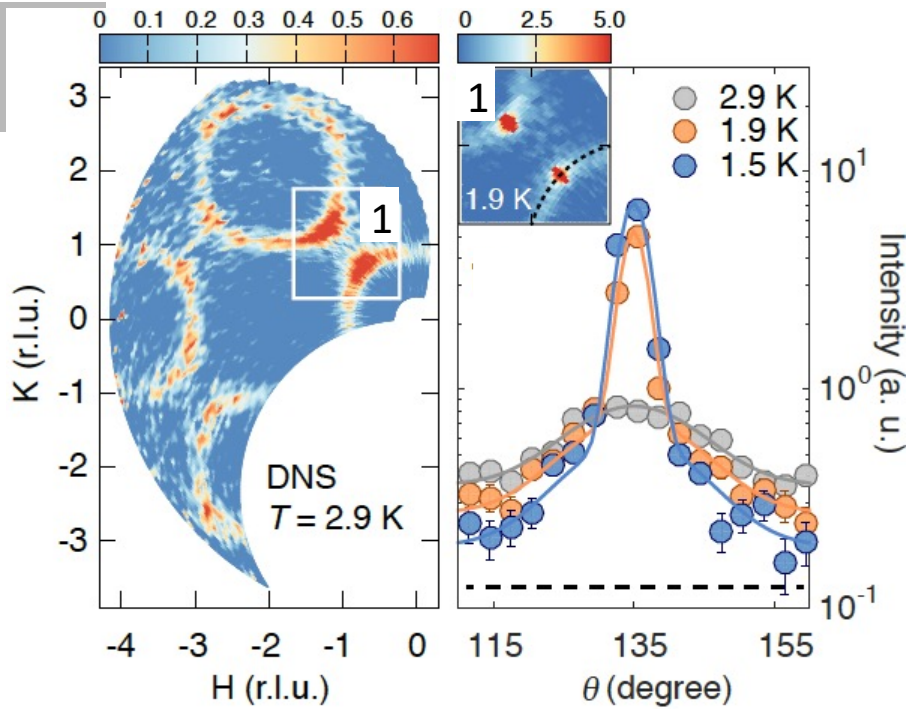
J_1 FM, J_2 AF

small AF $J_3 = -0.1J_1$

induces LRO $k = \frac{3}{4} \frac{3}{4} 0$



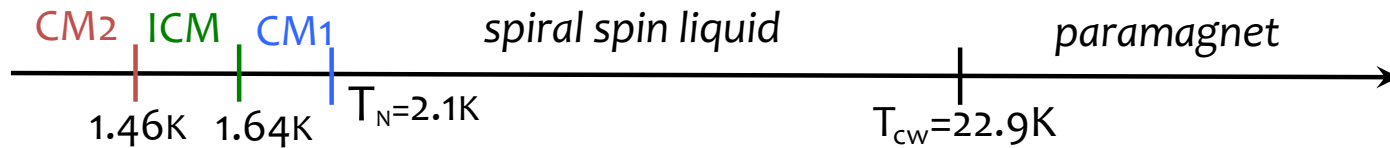
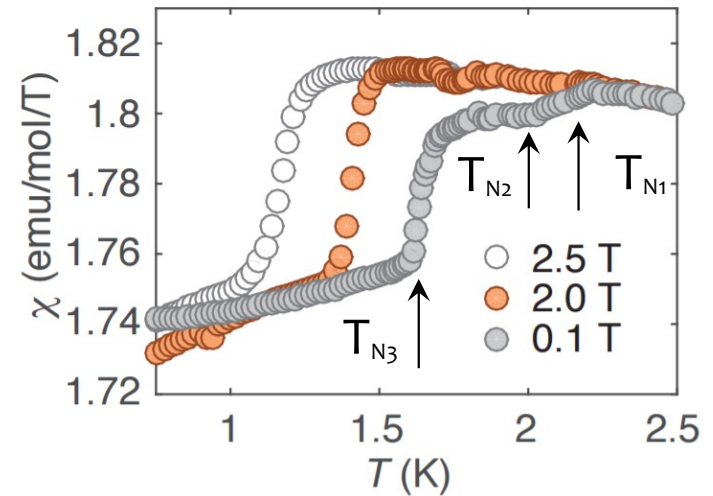
AF spinel MnSc₂S₄ and multistep ordering at H=0



$$T < T_N \quad k = \frac{3}{4} \frac{3}{4} 0$$

$$T_2 < T < T_1 \quad k = \frac{3}{4} \pm \delta \frac{3}{4} \mp \delta 0$$

$$T < T_2 \quad k = \frac{3}{4} \frac{3}{4} 0$$



AF spinel MnSc₂S₄ and multistep ordering

Neutron single crystal diffractometer ZEBRA

CM1 *amplitude modulated*

transversal to k

$$M_{1-10} \text{ for } k=(\frac{3}{4} \frac{3}{4} 0)$$

CM2 *helical structure* transversal to k

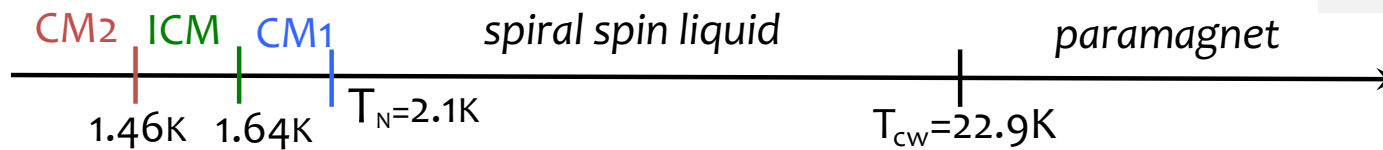
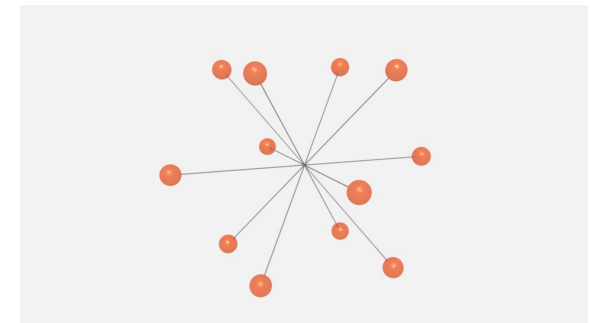
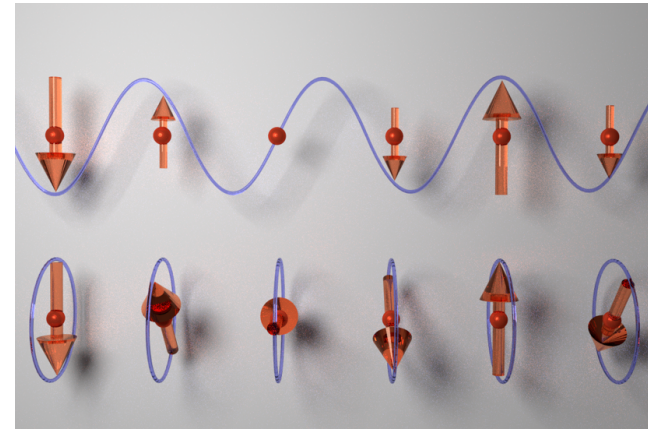
$$M_{1-10} \quad M_{001} \quad \text{for } k=(\frac{3}{4} \frac{3}{4} 0)$$

Cubic, $Fd-3m$, Mn^{2+} in $8a$ ($1/8 \ 1/8 \ 1/8$)

$k=(\frac{3}{4} \ \frac{3}{4} \ 0)$ the star has 12 arms, $-k \neq k$

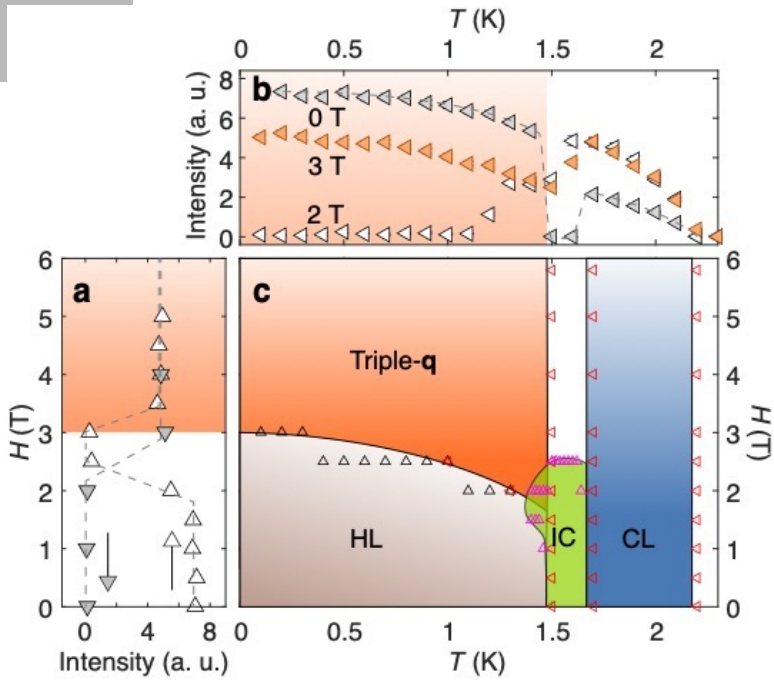
12 domain of single- k

4 domains of triple- k



AF spinel MnSc₂S₄ and magnetic field

Phase diagram H//001



Magnetic field favors states with highest magnetic susceptibility χ

Collinear domains with *spins* $\perp H$

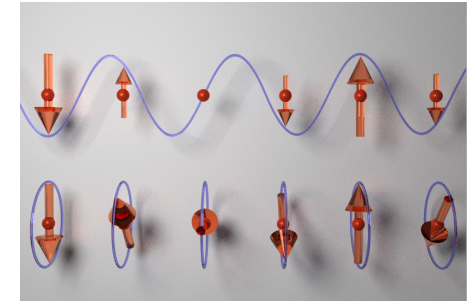
Helical domains with the *spin-plane* $\perp H$

Triple-k structure:

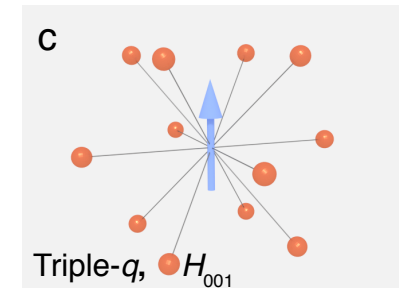
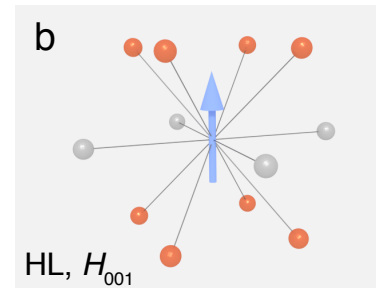
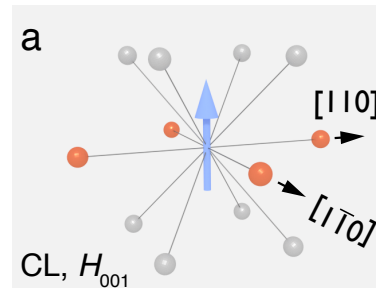
$$\Sigma (k_1 + k_2 + k_3) = 0$$

$$k_1 = [\frac{3}{4} \frac{3}{4} 0] \quad k_2 = [-\frac{3}{4} 0 -\frac{3}{4}] \quad k_3 = [0 -\frac{3}{4} \frac{3}{4}]$$

and $-k_1 -k_2 -k_3$

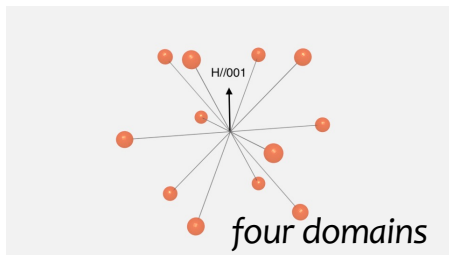
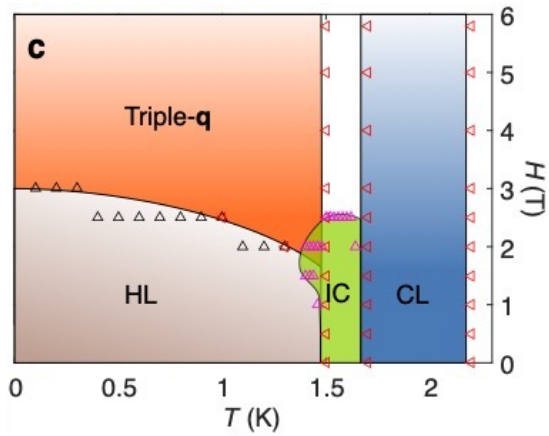


Expected magnetic reflections

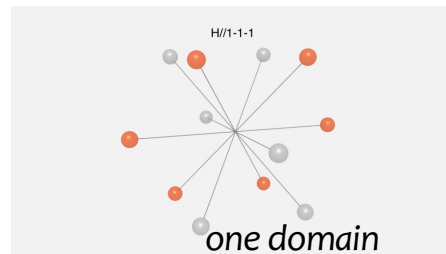
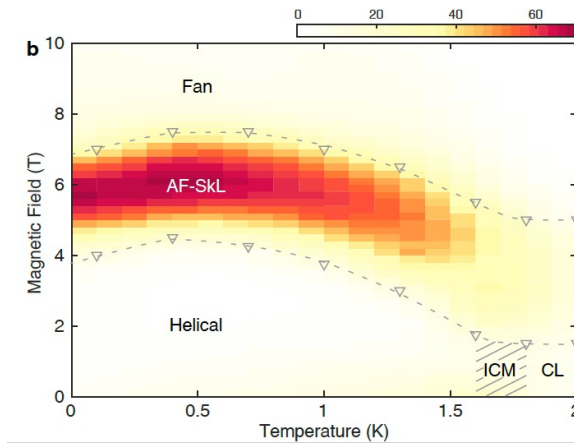


AF spinel MnSc₂S₄ and experimental (H,T) phase diagrams

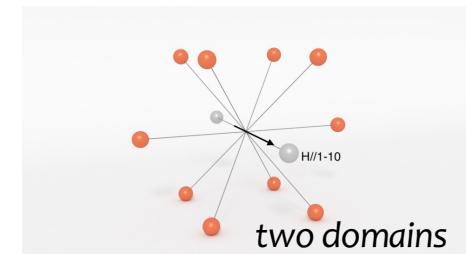
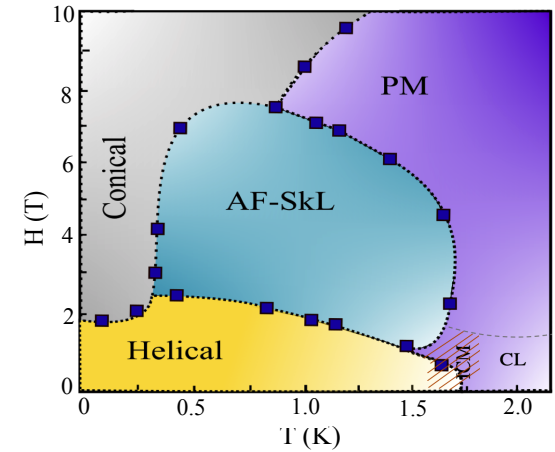
H // 001
ZEBRA



H // 111
D23

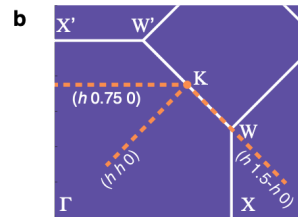
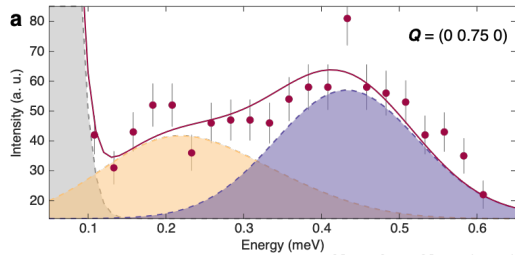


H // 110
ZEBRA

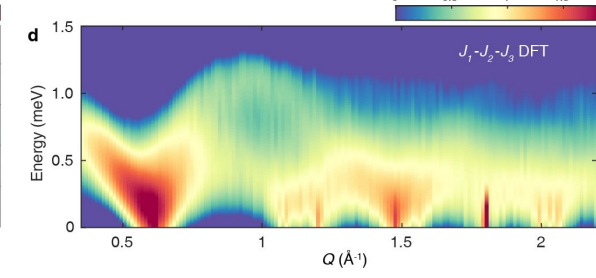
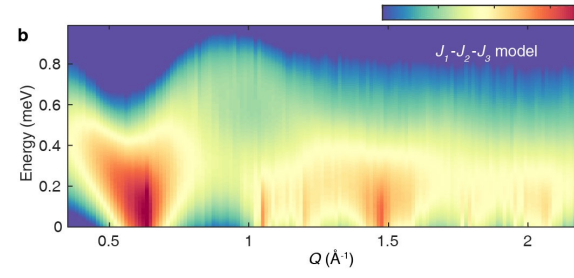
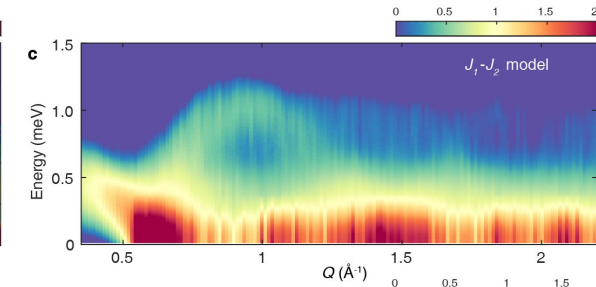
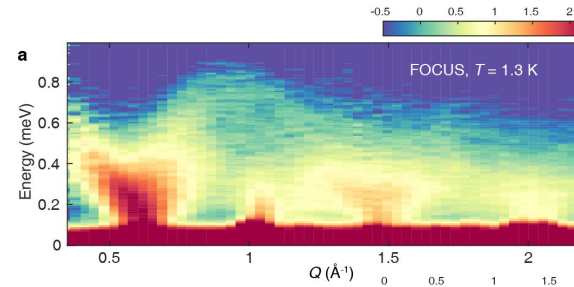
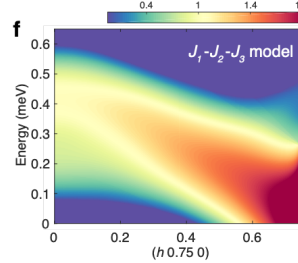
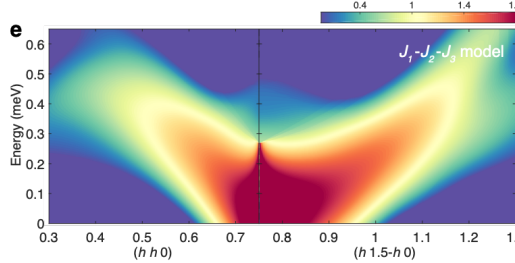
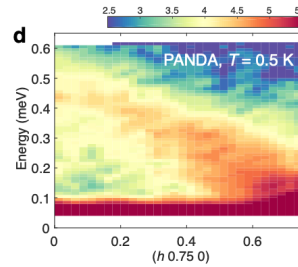
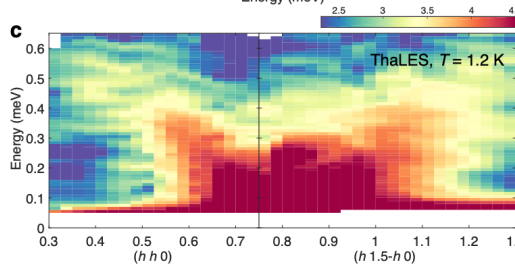


AF spinel MnSc₂S₄: INS

$$\mathcal{H} = \sum_{ij} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j + 3 J_{\parallel} \sum_{ij \in NN} (\mathbf{S}_i \cdot \hat{r}_{ij}) (\mathbf{S}_j \cdot \hat{r}_{ij}) + A_4 \sum_{i, \alpha=x,y,z} (S_i^\alpha)^4 - g \mu_B \mu_0 \mathbf{H} \sum_i \mathbf{S}_i$$



$J_1 = -0.31, J_2 = 0.46, J_3 = 0.087 \text{ K}$



AF spinel MnSc₂S₄: Monte-Carlo study

$$\mathcal{H} = \sum_{ij} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j + 3 J_{\parallel} \sum_{ij \in NN} (\mathbf{S}_i \cdot \hat{r}_{ij}) (\mathbf{S}_j \cdot \hat{r}_{ij}) + A_4 \sum_{i, \alpha=x,y,z} (S_i^\alpha)^4 - g \mu_B \mu_0 \mathbf{H} \sum_i \mathbf{S}_i$$

$$J_1 = -0.31, J_2 = 0.46, J_3 = 0.087 \text{ K}$$

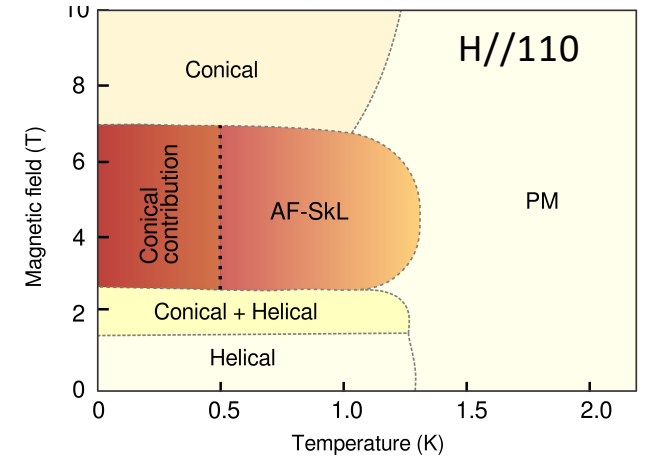
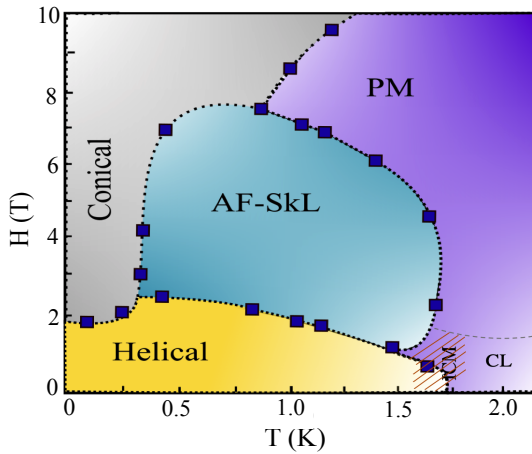
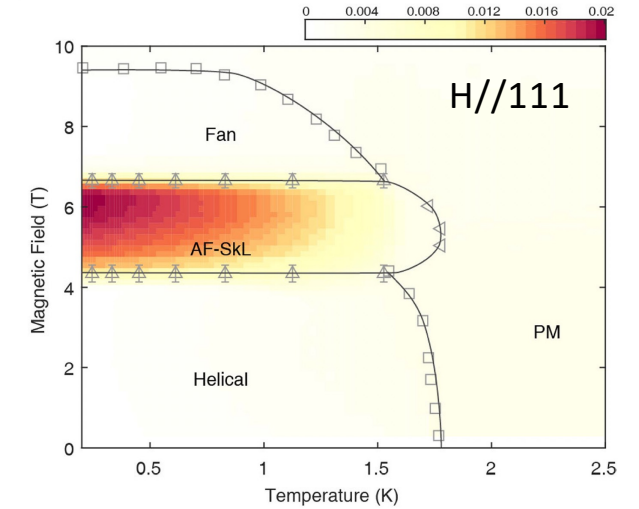
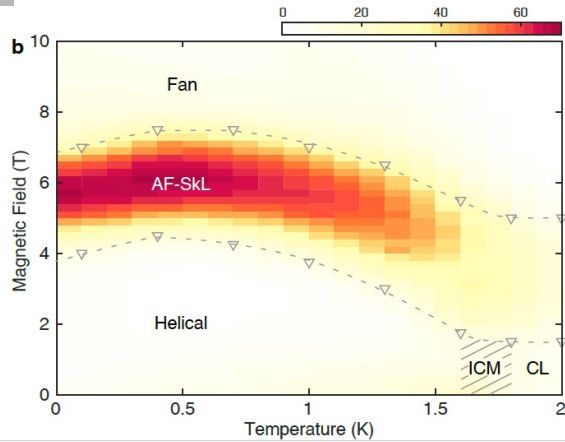
$$J_{\parallel} = -0.01 \text{ K}, A_4 = 0.0016 \text{ K}$$

Total scalar spin chirality

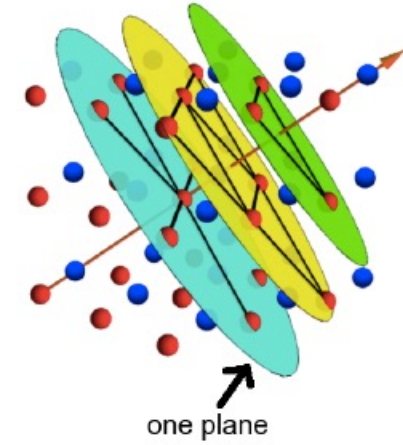
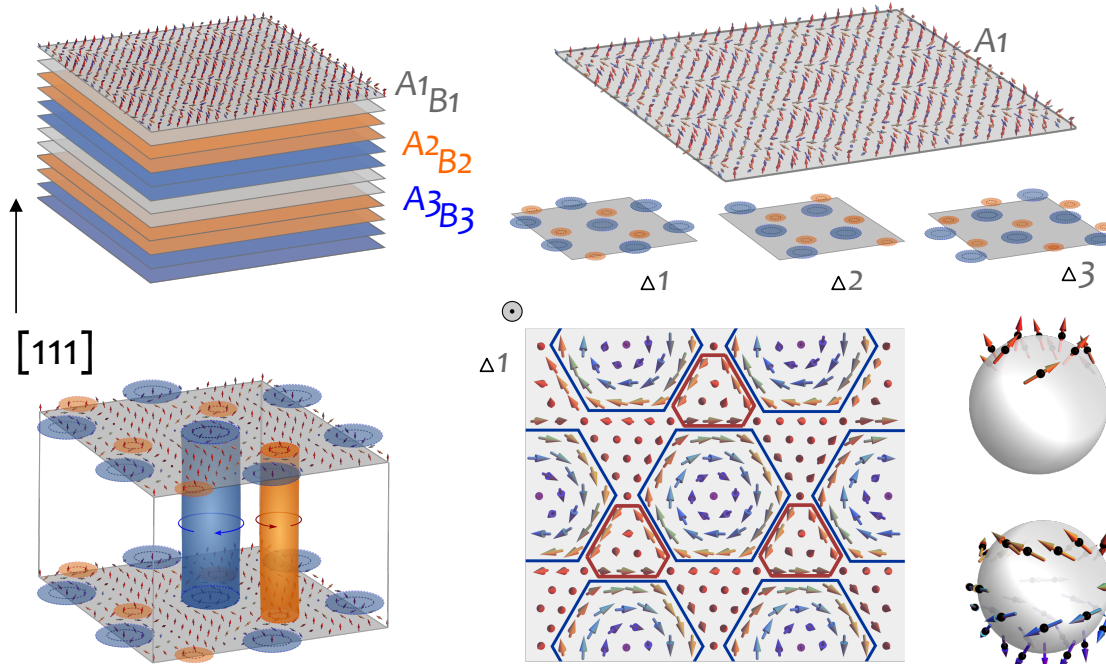
$$\chi_{tot} = \left\langle \frac{1}{N} \sum_n \chi_n \right\rangle$$

$$\chi_n = \mathbf{S}_i \cdot (\mathbf{S}_j \times \mathbf{S}_k)$$

fractional AF skyrmion lattice



AF spinel MnSc₂S₄: AF SkL description



A and B – two fcc sublattices ($t = 1/4 \ 1/4 \ 1/4$)
 A1, A2, A3 – triangular layers ($t = 1/2 \ 1/2 \ 0$)
 $\Delta 1, \Delta 2, \Delta 3$ – sublayer with AF SkL

$$\mathbf{S}_r = \mathbf{S}_r^\perp + \mathbf{S}_r^{[111]}$$

$$\mathbf{S}_r^\perp = \frac{1}{n_s} A^\perp \sum_{i=1}^3 \sin(\mathbf{Q}_i \cdot \mathbf{r} + \theta_i) \mathbf{e}_i$$

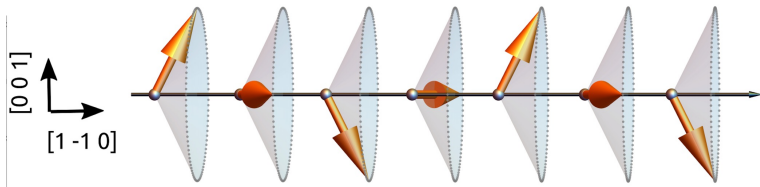
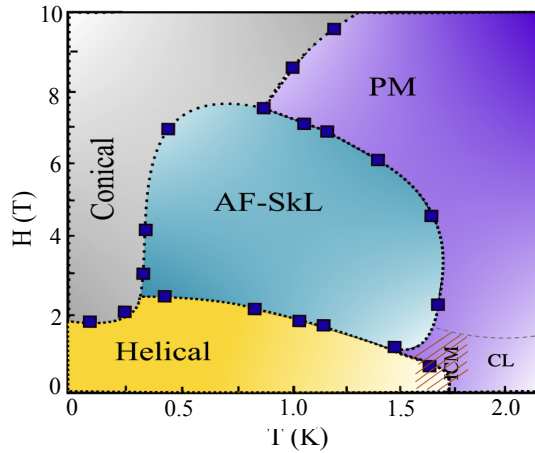
$$\mathbf{S}_r^{[111]} = \frac{1}{n_s} [M^{[111]} + A^{[111]} \sum_{i=1}^3 \cos(\mathbf{Q}_i \cdot \mathbf{r} + \phi_i)] \mathbf{e}^{[111]}$$

distance between cores $\sim 55 \text{ \AA}$

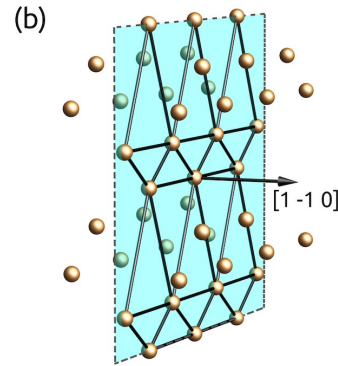
Bloch type

AF spinel MnSc₂S₄: locking of AF SkL axis

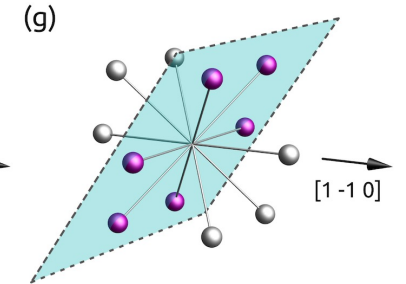
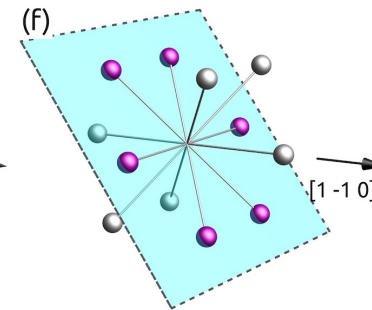
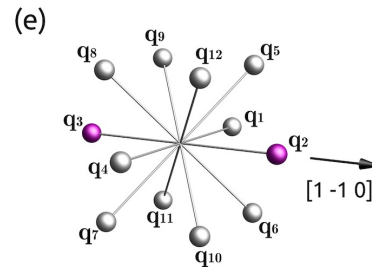
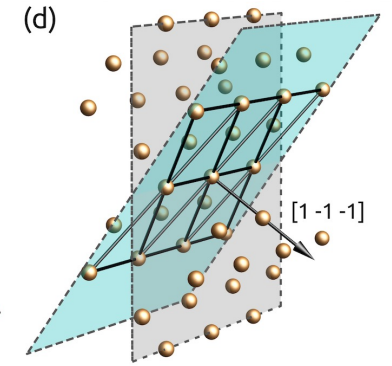
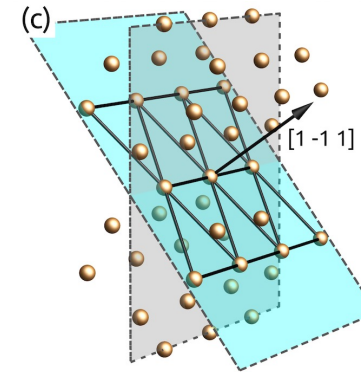
H // 110



Conical phase



Two domains of AF SkL phase



MnSc₂S₄ and other SkL

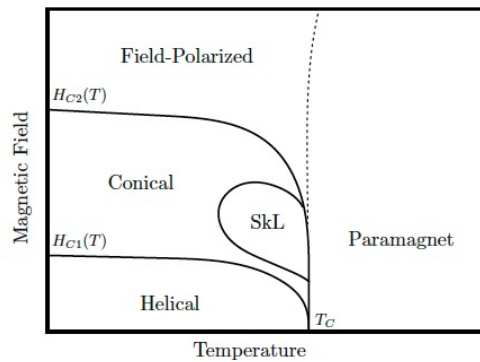
MnSi, Cu₂OSeO₃ B20-family

soft chiral ferromagnets

magnetic helix ~190 Å, skyrmions of the Bloch type

P_{2,3} non-centrosymmetric

FM interaction > DM interaction > anisotropy



GaV₄X₈ (X=S, Se) F-43m – R3m T_s = 42 K four polar domains

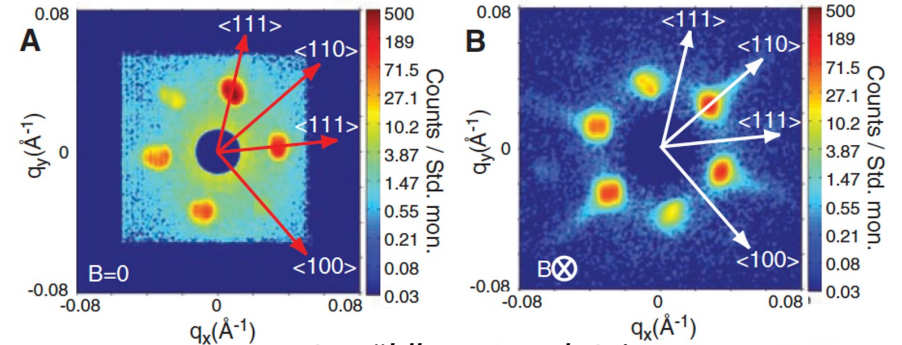
orientation confinement of SkL to the anisotropy axis

skyrmions of the Neel type distort and core displaces

SkL region is larger than in the B20 family

I. Kezsmarki et al. Nature Physics 14, 1116 2015

A. O. Leonov, I. Kezsmarki, PhysRevB. 96, 214413 2017



S. Mühlbauer et al. Science 323, 915 2009
SANS experiment

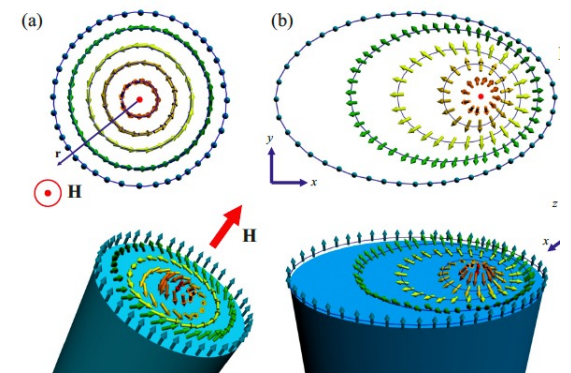
MnSc₂S₄

dominant AF coupling

centrosymmetric, as Gd₂PdSi₃

tiny anisotropy has strong effect

fractional Sk, d_{cores} ~ 55 Å



Main collaborators

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S. Gao et al. Nature 583, 37 2020

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