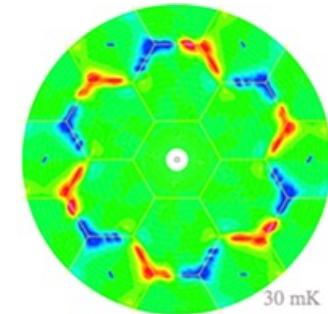




Chiral spin liquid ground state in $\text{YBaCo}_3\text{FeO}_7$

PHYSICAL REVIEW X 12, 021029 (2022)



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The system $\text{YBa}\text{Co}_3\text{FeO}_7$

$^{2+} \text{Co}^{3+}$
 $S=3/2, 5/2$

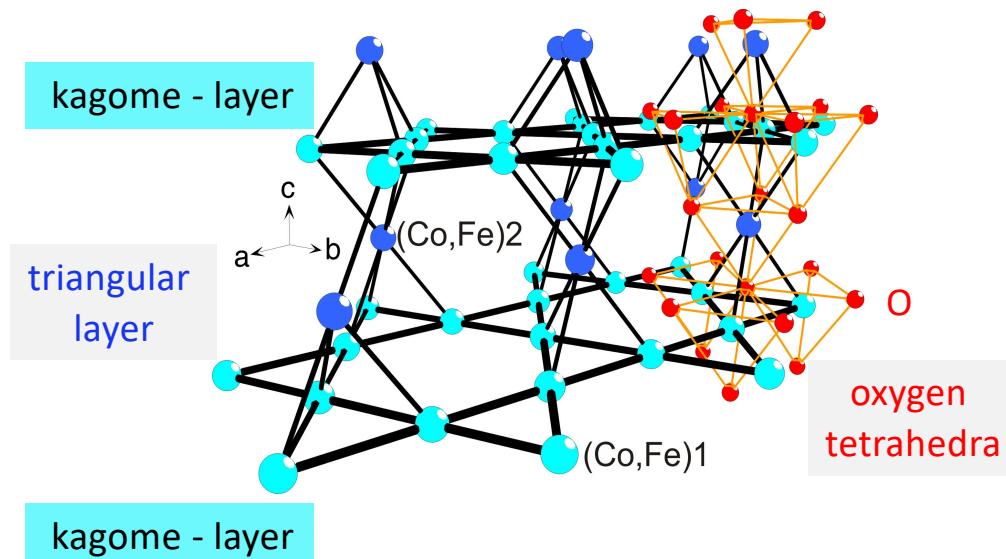
mineral $\text{SbNaBe}_4\text{O}_7$
structure Aminoff 1933, Pauling 1935

Structure

layered kagome-system

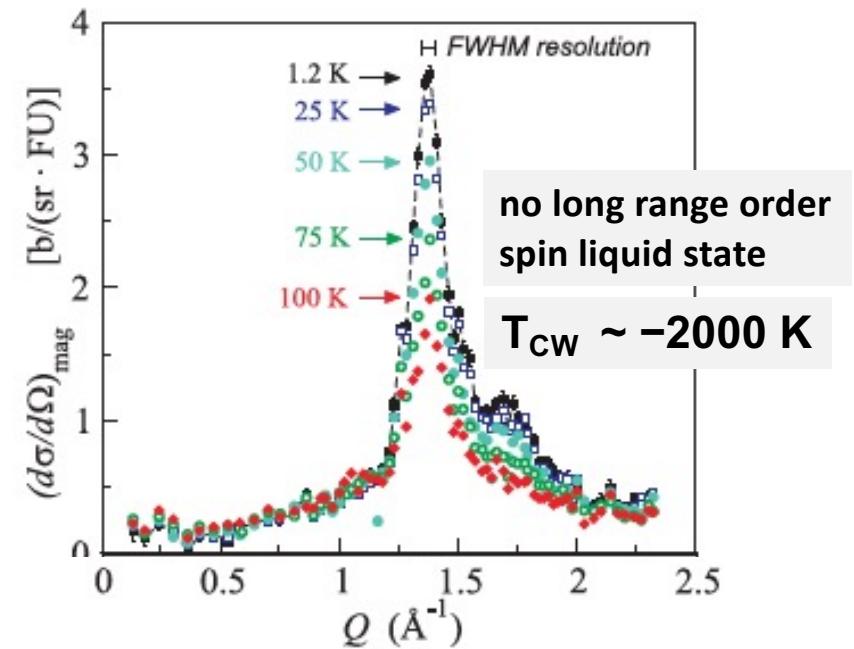
Swedenborgite structure

hexagonal $P6_3mc$



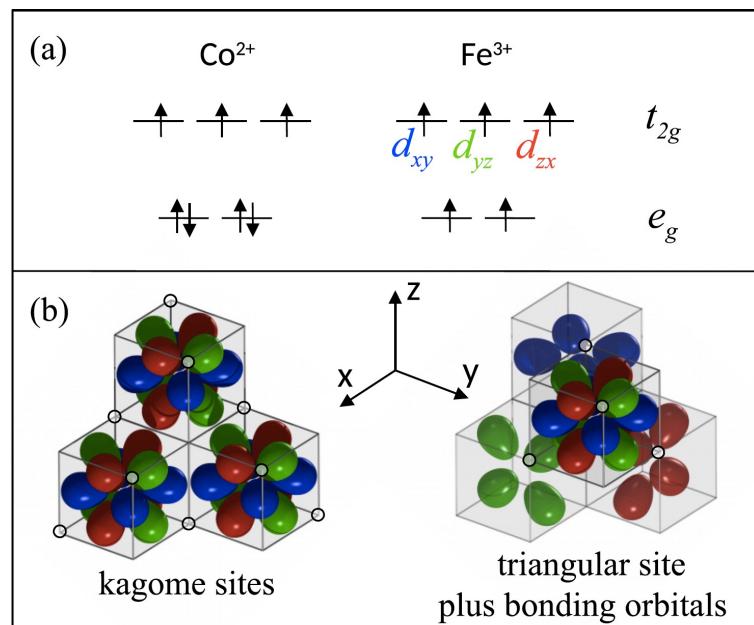
Spin correlations

polarized diffuse magnetic neutron scattering from powder



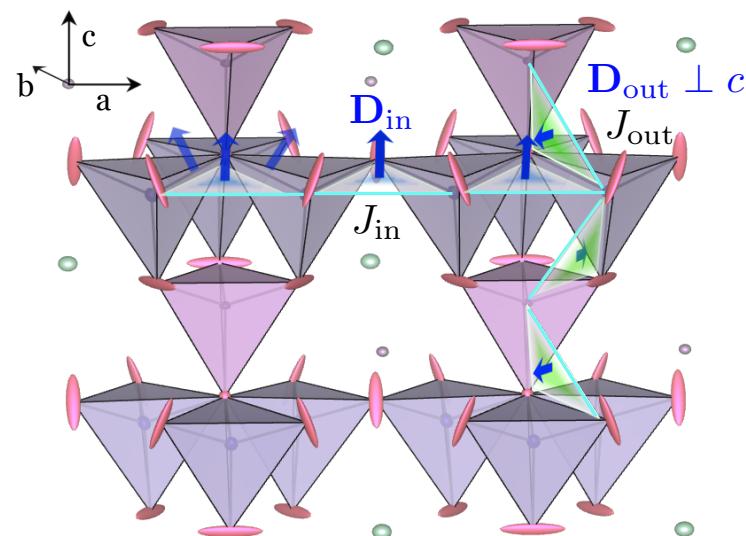
Valldor et al., Phys. Rev. B 84, 224426 (2011).

Exchange interactions



Super-exchange
between half-filled t_{2g} orbitals

Reim et al., Phys. Rev. B 97, 144402 (2018)



Dzyaloshinskii-Moriya interactions

$\mathbf{D} \Rightarrow$ vector chirality $\mathbf{C} = \langle \mathbf{S} \times \mathbf{S}' \rangle \neq 0$
 \mathbf{D}_{in} ab -plane anisotropy

3

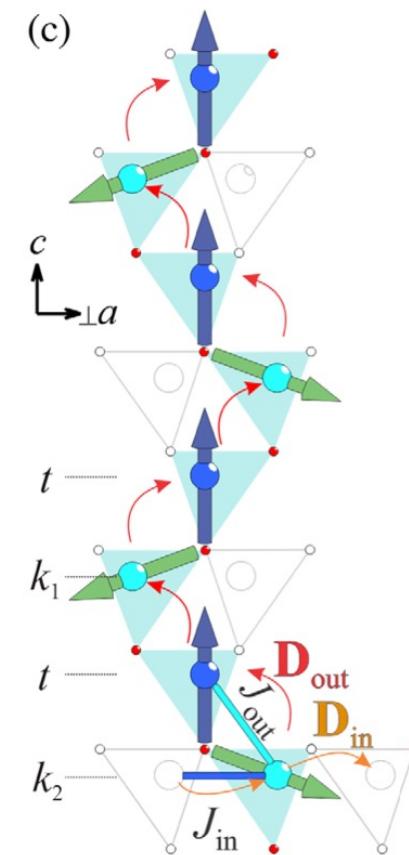
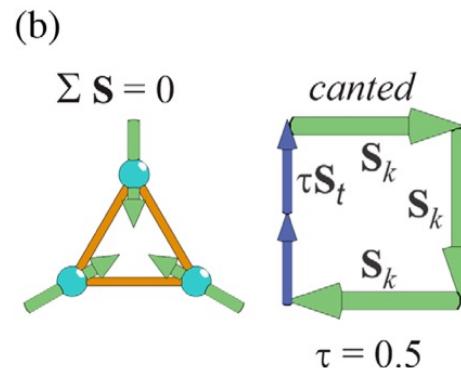
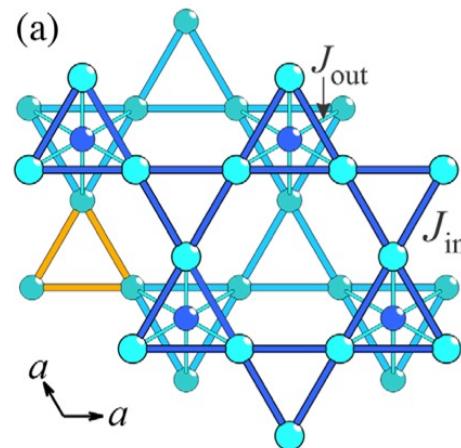
3/13

Frustration and sum rules

Manuel et al., Phys. Rev. Lett. 103, 037202 (2009).

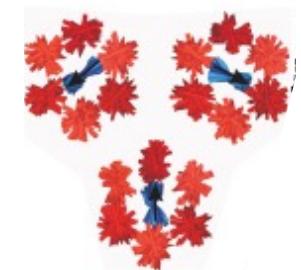
Possible Chiral Spin-Liquid Phase in Noncentrosymmetric $RBaCo_4O_7$

Khalyavin et al., Phys. Rev. B 85, 220401(R) (2012).



$$MCS$$

$$J_{in} = 2J_{out}$$



ferro-chains on t -sites

D_{out}
 t -site anisotropy $\parallel c$

D_{in}
 k -site anisotropy
towards ab -plane

weak anisotropies

XYZ-polarisation analysis

of diffuse magnetic neutron scattering from single crystals

*W. Schweika, J. Phys.: Conf. Ser. **211**, 012026 (2010)*

Powder samples



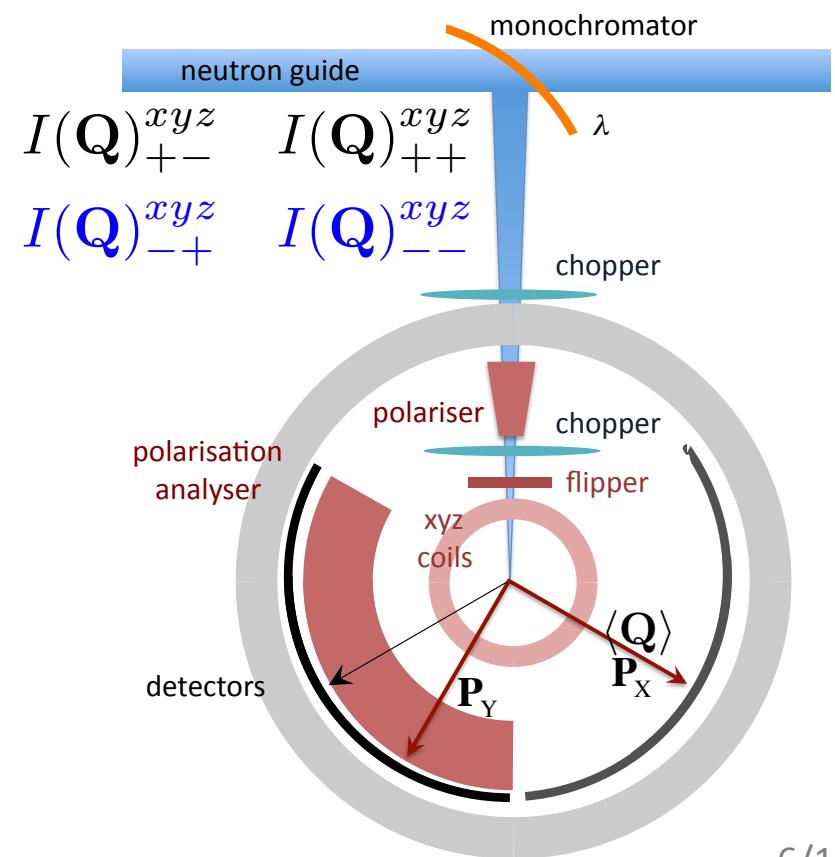
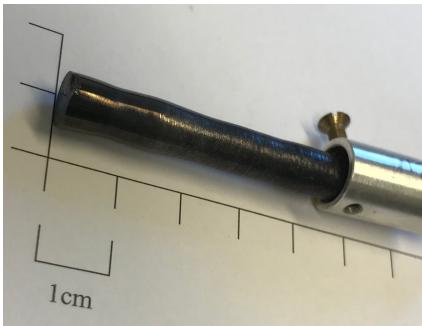
Otto Schärpf †13.6.2019

XYZ method 1993

D7 - ILL

XYZ-polarisation analysis of diffuse magnetic neutron scattering from single crystals

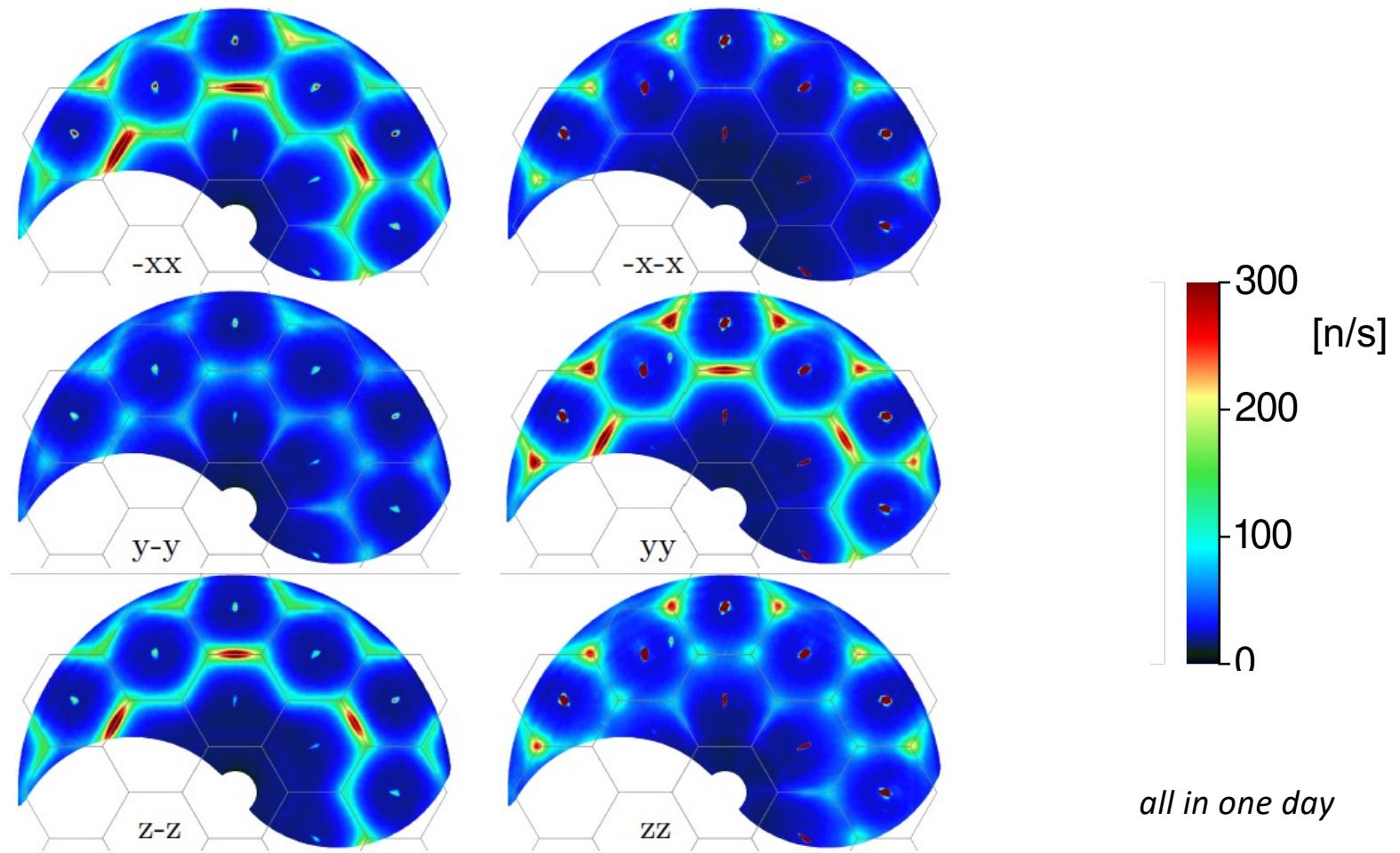
W. Schweika, J. Phys.: Conf. Ser. **211**, 012026 (2010)



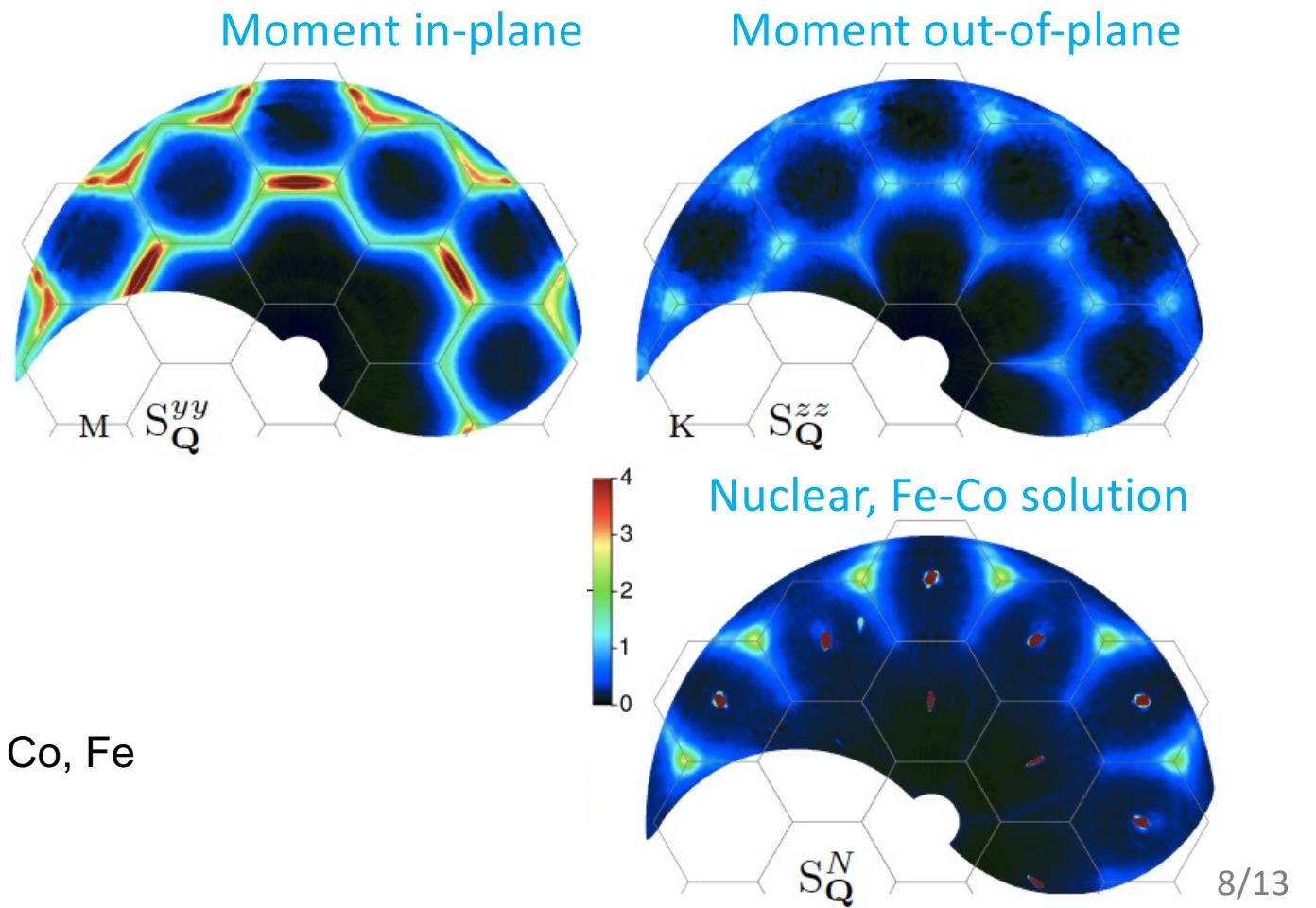
XYZ-polarisation analysis

of diffuse magnetic neutron scattering from single crystals

W. Schweika, J. Phys.: Conf. Ser. **211**, 012026 (2010)



Separation of diffuse nuclear and magnetic scattering

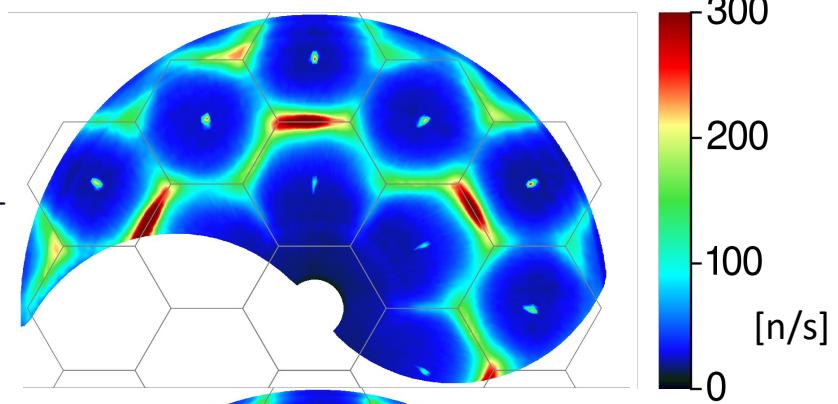


=> $S(Q)$
include average form factor Co, Fe
scale to $S=1$

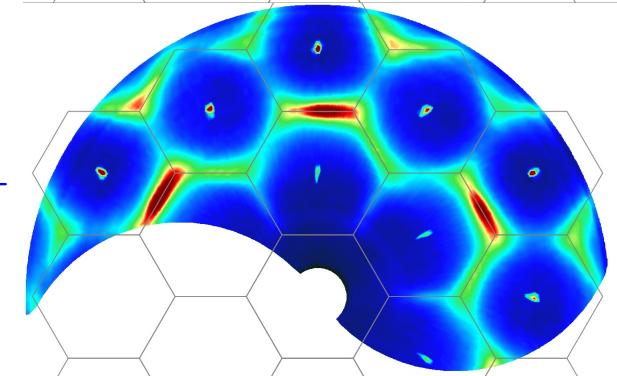
polarization reversal

$$\sigma_Q^{\text{chiral}} = i\mathbf{P} \cdot (\mathbf{M}_{-\mathbf{Q}}^\perp \times \mathbf{M}_\mathbf{Q}^\perp)$$

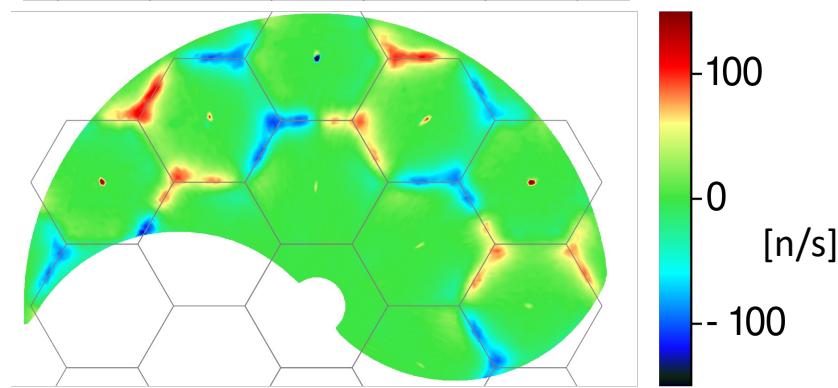
$$I(\mathbf{Q})_{+-}^x$$



$$I(\mathbf{Q})_{-+}^x$$



Difference: chiral scattering



Magnetic chiral diffuse scattering

$$\sigma_{\mathbf{Q}}^{\text{chiral}} = i\mathbf{P} \cdot (\mathbf{M}_{-\mathbf{Q}}^\perp \times \mathbf{M}_{\mathbf{Q}}^\perp)$$

Fourier analysis => vector chirality $\mathbf{C} = \mathbf{S}_{\mathbf{R}} \times \mathbf{S}'_{\mathbf{R}'}$ $S^y S^z$ ($\perp \mathbf{Q}$)

$$S_{\mathbf{Q}}^{yz} = i\mathbf{P} \cdot \mathbf{C}(\mathbf{Q})$$

propagation

$$= \sum_{\mathbf{r}} \mathbf{e}_{\mathbf{Q}} \cdot \begin{pmatrix} G_{x,\mathbf{r}}^{yz} \mathbf{e}_{\mathbf{r}\parallel} \\ G_{y,\mathbf{r}}^{yz} \mathbf{e}_{\mathbf{r}\perp^y} \\ G_{z,\mathbf{r}}^{yz} \mathbf{e}_{\mathbf{r}\perp^z} \end{pmatrix} \sin(\mathbf{Q} \cdot \mathbf{r})$$

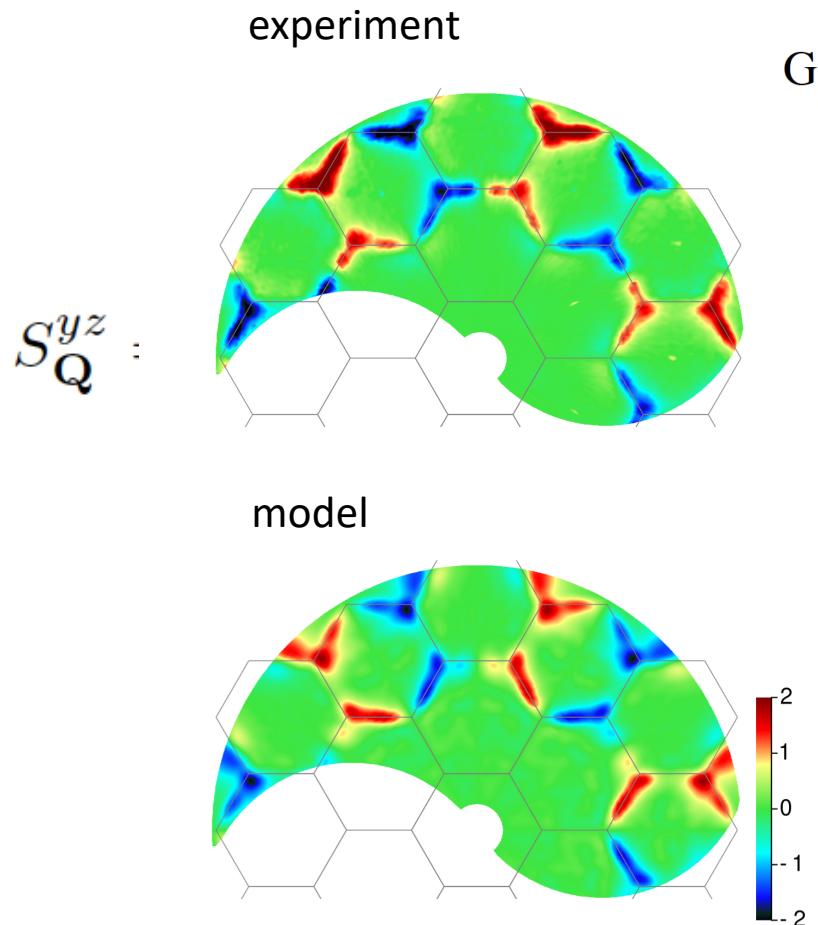
x	Helix
y	Cycloid
z	Cycloid

$$= \sum_{\mathbf{r}} \begin{pmatrix} G_{x,\mathbf{r}}^{yz} \mathbf{e}_{\mathbf{Q}} \cdot \mathbf{e}_{\mathbf{r}} \\ G_{y,\mathbf{r}}^{yz} \|\mathbf{e}_{\mathbf{Q}} \times \mathbf{e}_{\mathbf{r}}\| \\ 0 \end{pmatrix} \sin(\mathbf{Q} \cdot \mathbf{r})$$

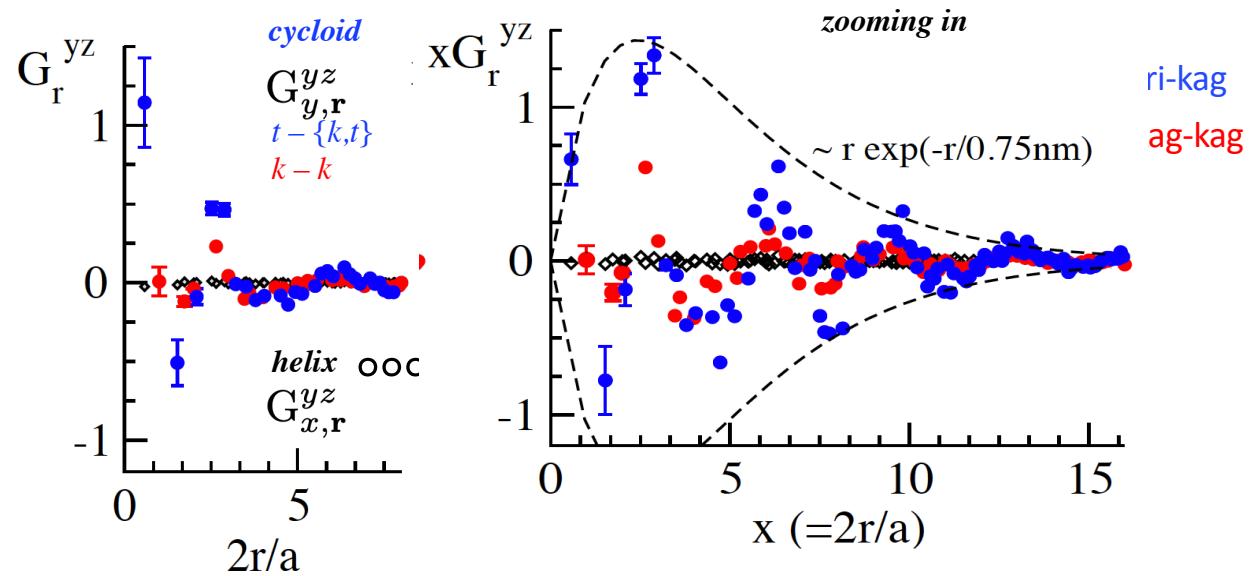
$\parallel \mathbf{Q}$	Helix
$\perp \mathbf{Q}$	Cycloid
$\perp \mathbf{Q}$	(not measured in (hk0))

anti-symmetry

Magnetic chiral diffuse scattering



Fourier analysis

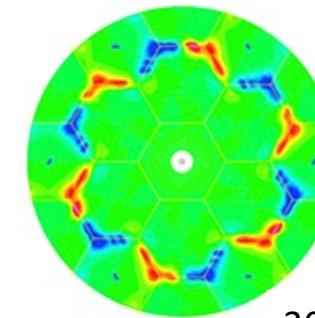
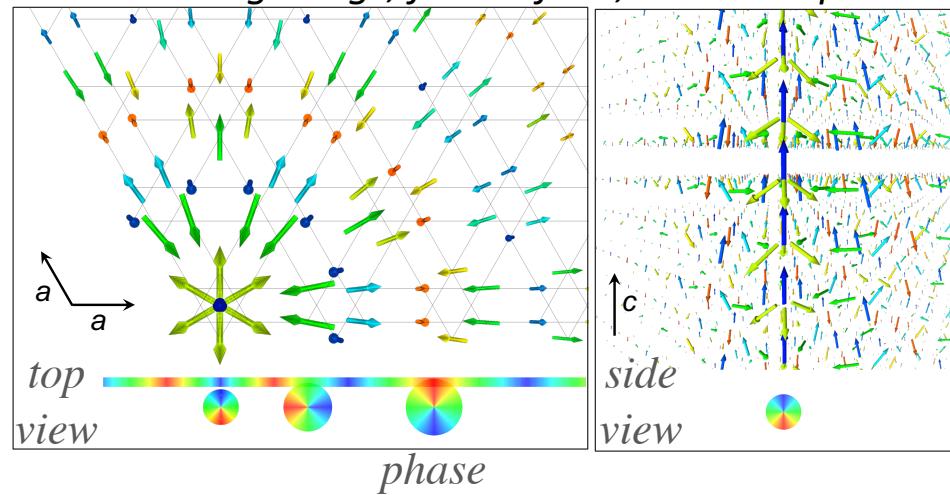


3f-cycloids emanating from t -site
radial cycloidal wave

Structure and nature

3f - cycloidal structures emanate from t -sites
suggesting zero dimensional objects “lumps” stable at lowest T
requiring sufficiently weak anisotropies

similar to *hedge hogs, flower field, water drops on sea surface*



30mK

Could it be
a quantum spin liquid state?
Skyrmions?

Theoretical analysis

Ulrich Rößler (Skyrmion ground states, Nature 2006)

Standard classical theory of magnetic order shows the existence of relevant Lifshitz invariants for the non-centro symmetric and polar structure. Coupling of higher order terms due to chirality may twist the spin structure into an inextricable structure of **short-range entities we term lumps**.

Existence of Lifshitz-type invariants for C_{6v}

$$\begin{array}{ll} \text{cycloids in basal plane } \perp c & g_{lm}(l_x\partial_x m_z - m_z\partial_x l_x + l_y\partial_y m_z - m_z\partial_y l_y) \\ \text{modulation along the } c \text{ axis} & f_{lm}(l_x\partial_z m_x - m_x\partial_z l_x) \end{array}$$

As in the case of skyrmions in chiral helimagnets, the free energy density in the cores of these lumps is reduced compared to a single spiral, by activating several Lifshitz-terms. **This stabilizes these static solitonic objects and generates an inhomogeneous state that constitutes a classical ground state.**

The present scenario with antisymmetric exchange acting as a frustrating gauge background that stabilizes local spin lumps, is similar to the avoided phase transition in coupled gauge- and matter-fields for sub-nuclear particles, and **may emerge in other non-centro symmetric and highly frustrated systems.**