

Measurement of the Fierz interference term with PERKEO III

Karina Bernert

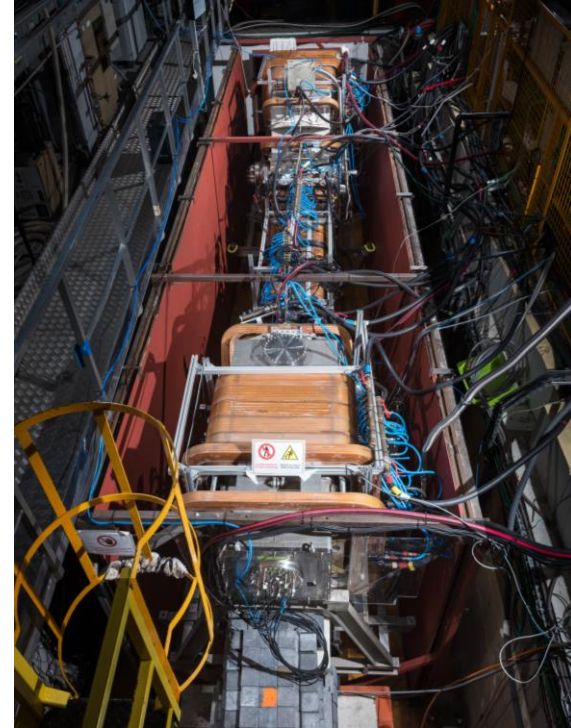
Technical University of Munich

Physics Department

Particle Physics at Low Energy

Joint ESS ILL User Meeting

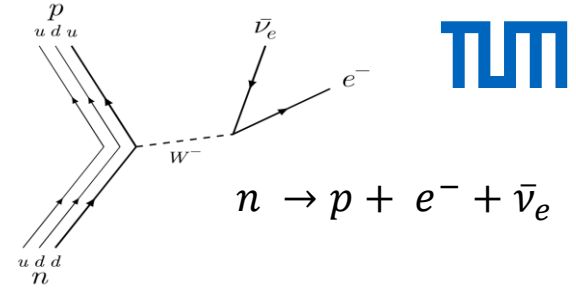
Lund, 7th of October 2022



PERKEO III at PF1b
©2019 Laurent Thion <ecliptique.com>

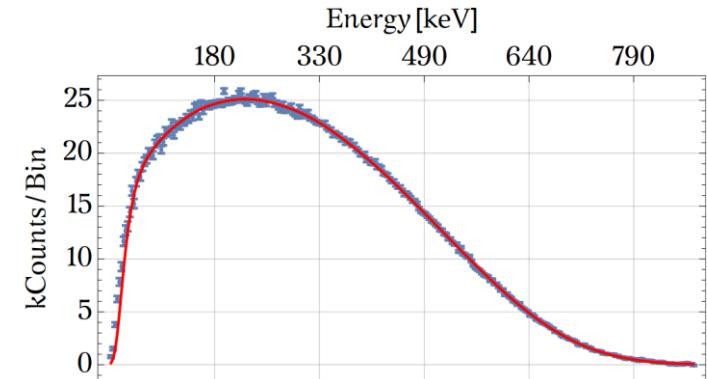
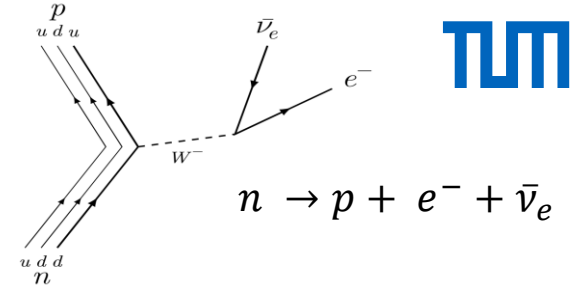
Search for BSM physics with neutron decay

Testing the Standard Model with
high precision spectroscopy



Search for BSM physics with neutron decay

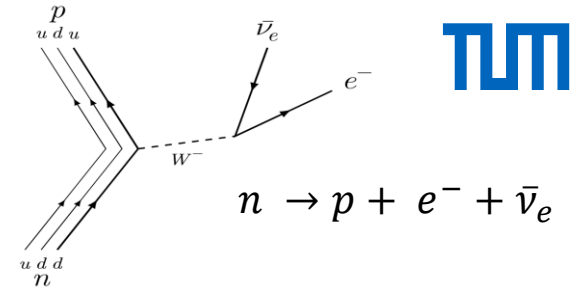
Testing the Standard Model with
high precision spectroscopy



PhD Thesis C. Roick 2018

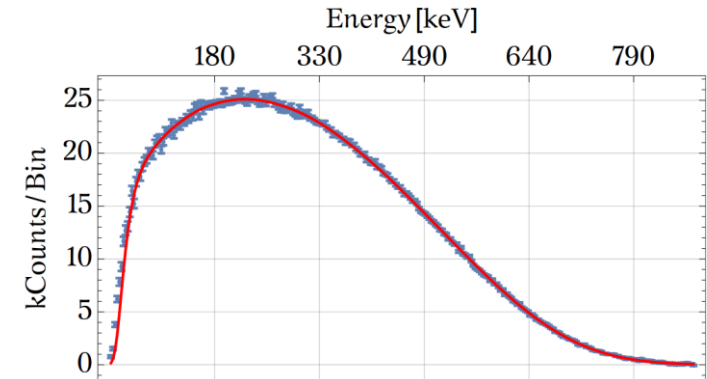
Search for BSM physics with neutron decay

Testing the Standard Model with high precision spectroscopy



Differential decay width can be described in terms of several correlation coefficients: [JTW57]

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$



[JTW57]: Jackson, Treiman, Wyld (1957). Possible Tests of Time Reversal Invariance in Beta Decay

Search for BSM physics with neutron decay

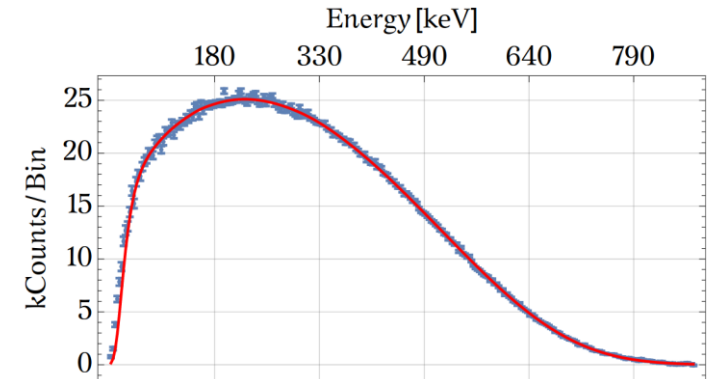
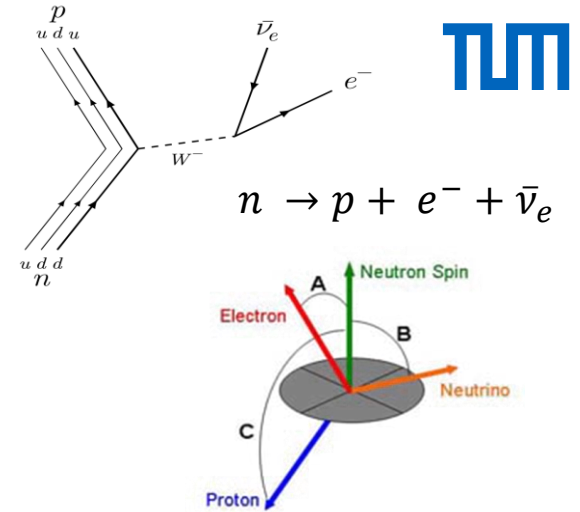
Testing the Standard Model with high precision spectroscopy

Differential decay width can be described in terms of several correlation coefficients: [JTW57]

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

| | |
|---|---------------------------------------|
| a | Electron Neutrino Angular Correlation |
| b | Fierz Interference Term |
| A | Beta Asymmetry |
| B | Neutrino Asymmetry |
| C | Proton Asymmetry |
| D | Triple Correlation Coefficient |

Coefficients are either 0 or depend on $\lambda = g_A/g_V$ within SM



PhD Thesis C. Roick 2018

[JTW57]: Jackson, Treiman, Wyld (1957). Possible Tests of Time Reversal Invariance in Beta Decay

Search for BSM physics with neutron decay

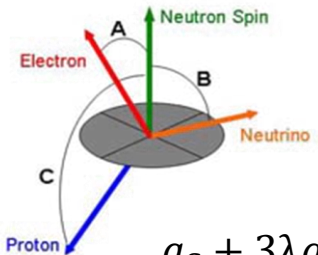
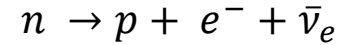
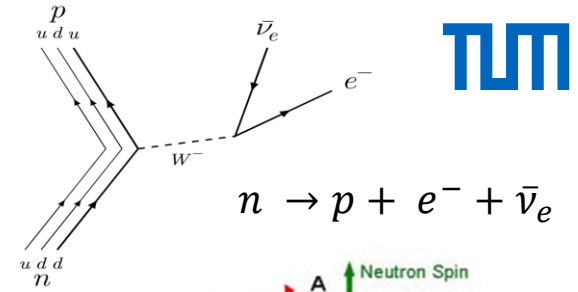
Testing the Standard Model with high precision spectroscopy

Differential decay width can be described in terms of several correlation coefficients: [JTW57]

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

| | |
|---|---------------------------------------|
| a | Electron Neutrino Angular Correlation |
| b | Fierz Interference Term |
| A | Beta Asymmetry |
| B | Neutrino Asymmetry |
| C | Proton Asymmetry |
| D | Triple Correlation Coefficient |

Coefficients are either 0 or depend on $\lambda = g_A/g_V$ within SM



$$b \cong \frac{g_S g_V + 3g_A g_T}{g_V^2 + g_S^2 + 3(g_A^2 + g_T^2)} \cong 2 \frac{g_S + 3\lambda g_T}{1 + 3\lambda^2}$$

- There are no scalar or tensor interactions in SM

$$\rightarrow g_S = g_T = 0 \rightarrow b = 0$$

- But if we measure $b \neq 0$

$$\rightarrow g_S \neq 0 \text{ or } g_T \neq 0$$

→ New physics

[JTW57]: Jackson, Treiman, Wyld (1957). Possible Tests of Time Reversal Invariance in Beta Decay

Two ways to get the Fierz term b

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

Two ways to get the Fierz term b

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

1st Method: using polarized neutrons

Combined fit of b and electron asymmetry A

$$A_{\text{exp}}(E_e) = \frac{N^\uparrow(E_e) - N^\downarrow(E_e)}{N^\uparrow(E_e) + N^\downarrow(E_e)} = \frac{1}{2} A(\lambda) P_n \beta(E_e)$$

$$A'_{\text{exp}}(E_e) \rightarrow A_{\text{exp}}(E_e) \cdot \frac{1}{1 + b \frac{m_e}{E_e}}$$

Two ways to get the Fierz term b

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

1st Method: using polarized neutrons

Combined fit of b and electron asymmetry A

$$A_{\text{exp}}(E_e) = \frac{N^\uparrow(E_e) - N^\downarrow(E_e)}{N^\uparrow(E_e) + N^\downarrow(E_e)} = \frac{1}{2} A(\lambda) P_n \beta(E_e)$$

$$A'_{\text{exp}}(E_e) \rightarrow A_{\text{exp}}(E_e) \cdot \frac{1}{1 + b \frac{m_e}{E_e}}$$

Statistical sensitivity, background contributions, and detector resolution affect the parameter correlations between λ and b

→ Statistical uncertainty dominates!

Two ways to get the Fierz term b

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

1st Method: using polarized neutrons

Combined fit of b and electron asymmetry A

$$A_{\text{exp}}(E_e) = \frac{N^\uparrow(E_e) - N^\downarrow(E_e)}{N^\uparrow(E_e) + N^\downarrow(E_e)} = \frac{1}{2} A(\lambda) P_n \beta(E_e)$$

$$A'_{\text{exp}}(E_e) \rightarrow A_{\text{exp}}(E_e) \cdot \frac{1}{1 + b \frac{m_e}{E_e}}$$

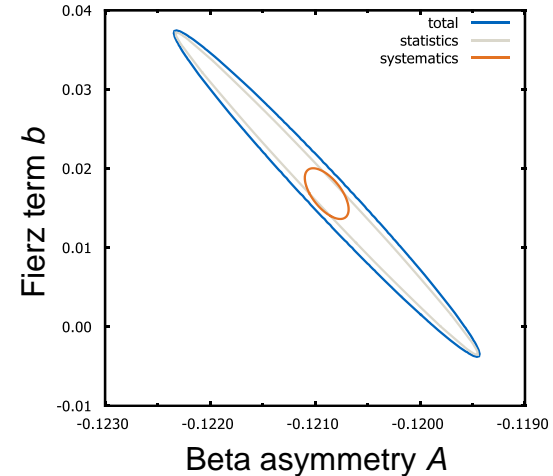
Statistical sensitivity, background contributions, and detector resolution affect the parameter correlations between λ and b

→ **Statistical uncertainty dominates!**

$$\text{PERKEO III}^1: b = 0.017(20)_{\text{stat}}(3)_{\text{sys}} = 0.017(21)$$

Current best value

from free neutron decay



¹ Saul et al., Phys. Rev. Lett. 125, 112501 (2020)

Two ways to get the Fierz term b

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

1st Method: using polarized neutrons

Combined fit of b and electron asymmetry A

$$A_{\text{exp}}(E_e) = \frac{N^\uparrow(E_e) - N^\downarrow(E_e)}{N^\uparrow(E_e) + N^\downarrow(E_e)} = \frac{1}{2} A(\lambda) P_n \beta(E_e)$$

$$A'_{\text{exp}}(E_e) \rightarrow A_{\text{exp}}(E_e) \cdot \frac{1}{1 + b \frac{m_e}{E_e}}$$

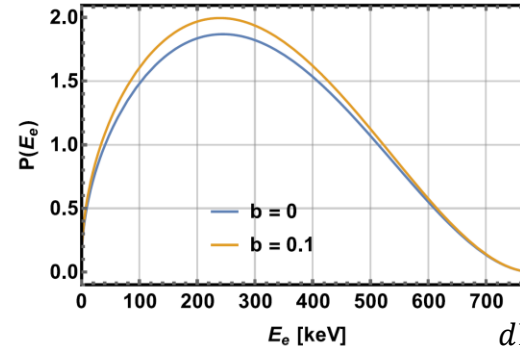
Statistical sensitivity, background contributions, and detector resolution affect the parameter correlations between λ and b

→ **Statistical uncertainty dominates!**

$$\text{PERKEO III}^1: b = 0.017(20)_{\text{stat}}(3)_{\text{sys}} = 0.017(21)$$

Current best value
from free neutron decay

2nd Method: using unpolarized neutrons



$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} \right)$$

¹ Saul et al., Phys. Rev. Lett. 125, 112501 (2020)

Two ways to get the Fierz term b

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

1st Method: using polarized neutrons

Combined fit of b and electron asymmetry A

$$A_{\text{exp}}(E_e) = \frac{N^\uparrow(E_e) - N^\downarrow(E_e)}{N^\uparrow(E_e) + N^\downarrow(E_e)} = \frac{1}{2} A(\lambda) P_n \beta(E_e)$$

$$A'_{\text{exp}}(E_e) \rightarrow A_{\text{exp}}(E_e) \cdot \frac{1}{1 + b \frac{m_e}{E_e}}$$

Statistical sensitivity, background contributions, and detector resolution affect the parameter correlations between λ and b

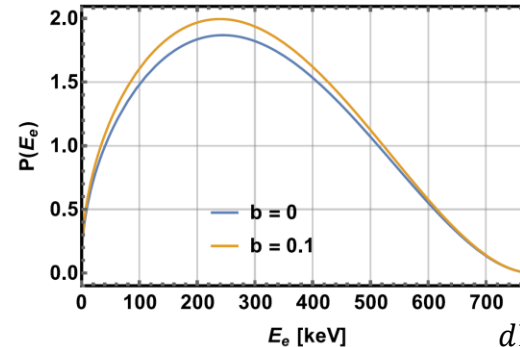
→ **Statistical uncertainty dominates!**

$$\text{PERKEO III}^1: b = 0.017(20)_{\text{stat}}(3)_{\text{sys}} = 0.017(21)$$

Current best value

from free neutron decay

2nd Method: using unpolarized neutrons



$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} \right)$$

First result from free neutron decay by UCNA²:

$$b = 0.067(0.005)_{\text{stat}} \begin{pmatrix} +0.090 \\ -0.061 \end{pmatrix}_{\text{sys}}$$

¹ Saul et al., Phys. Rev. Lett. 125, 112501 (2020)

² Hickerson et al., Phys. Rev. C 96 (2017)

Two ways to get the Fierz term b

$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} + a \frac{\vec{p}_e \cdot \vec{p}_\nu}{E_e E_\nu} + \hat{s}_n \cdot \left[A \frac{\vec{p}_e}{E_e} + B \frac{\vec{p}_\nu}{E_\nu} + D \frac{\vec{p}_e \times \vec{p}_\nu}{E_e E_\nu} \right] \right)$$

1st Method: using polarized neutrons

Combined fit of b and electron asymmetry A

$$A_{\text{exp}}(E_e) = \frac{N^\uparrow(E_e) - N^\downarrow(E_e)}{N^\uparrow(E_e) + N^\downarrow(E_e)} = \frac{1}{2} A(\lambda) P_n \beta(E_e)$$

$$A'_{\text{exp}}(E_e) \rightarrow A_{\text{exp}}(E_e) \cdot \frac{1}{1 + b \frac{m_e}{E_e}}$$

Statistical sensitivity, background contributions, and detector resolution affect the parameter correlations between λ and b

→ **Statistical uncertainty dominates!**

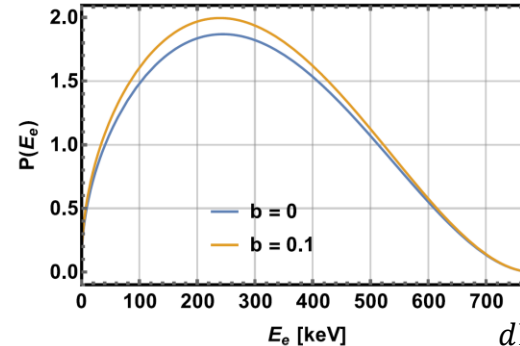
$$\text{PERKEO III}^1: b = 0.017(20)_{\text{stat}}(3)_{\text{sys}} = 0.017(21)$$

Current best value
from free neutron decay

¹ Saul et al., Phys. Rev. Lett. 125, 112501 (2020)

² Hickerson et al., Phys. Rev. C 96 (2017)

2nd Method: using unpolarized neutrons



$$d\Gamma_n \propto \rho(E_e) \left(1 + b \frac{m_e}{E_e} \right)$$

First result from free neutron decay by UCNA²:

$$b = 0.067(0.005)_{\text{stat}} \begin{pmatrix} +0.090 \\ -0.061 \end{pmatrix}_{\text{sys}}$$

→ **Systematic uncertainty dominates!**

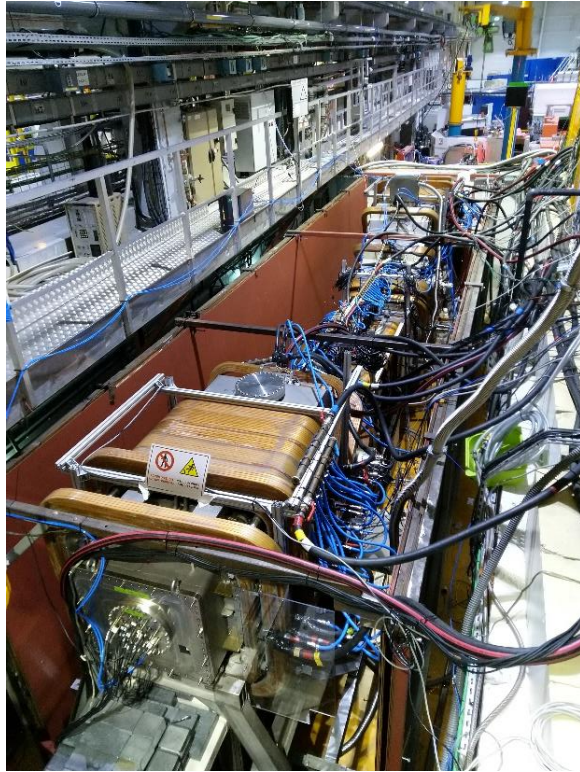
2019/2020 PERKEO III campaign at the ILL

- Measurement of Fierz term b using unpolarized neutrons (Method 2)
- Total uncertainty goal:
 $\Delta b \sim 5 \times 10^{-3}$
(Factor of four improvement to present limit using free neutron decay!)

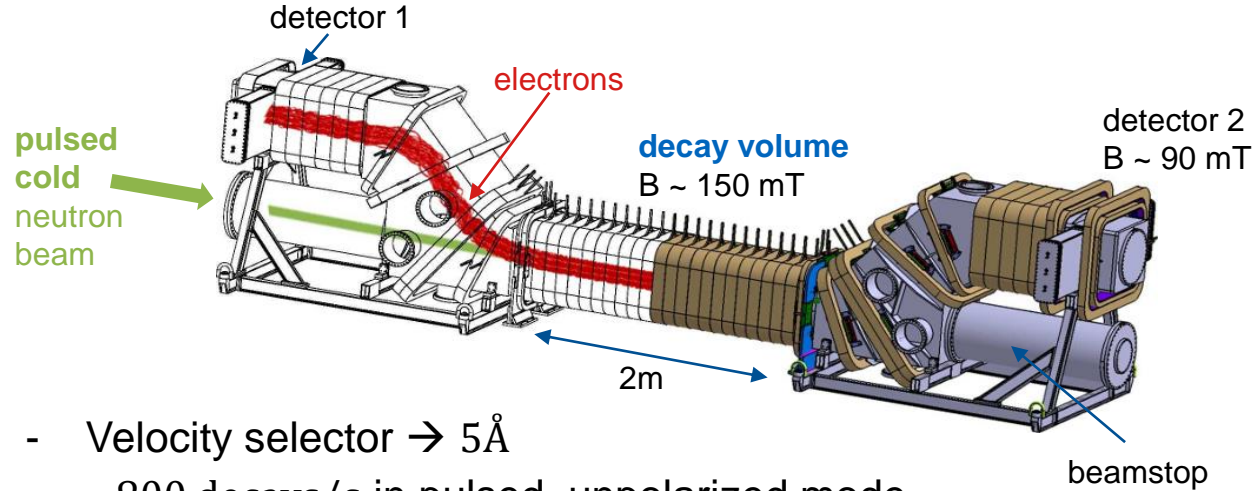


TU Munich, ILL, TU Vienna, University of Heidelberg

PERKEO III



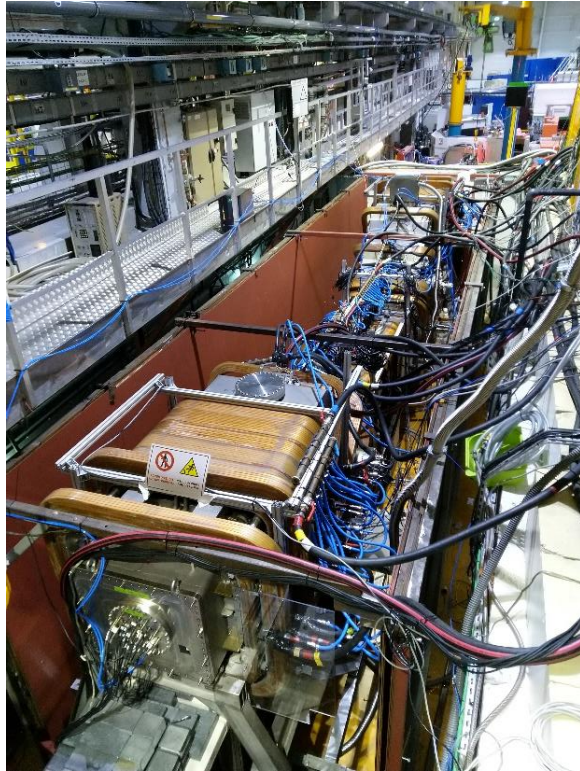
Originally built by University of Heidelberg,
now operated by TUM, TU Vienna, HD & ILL.



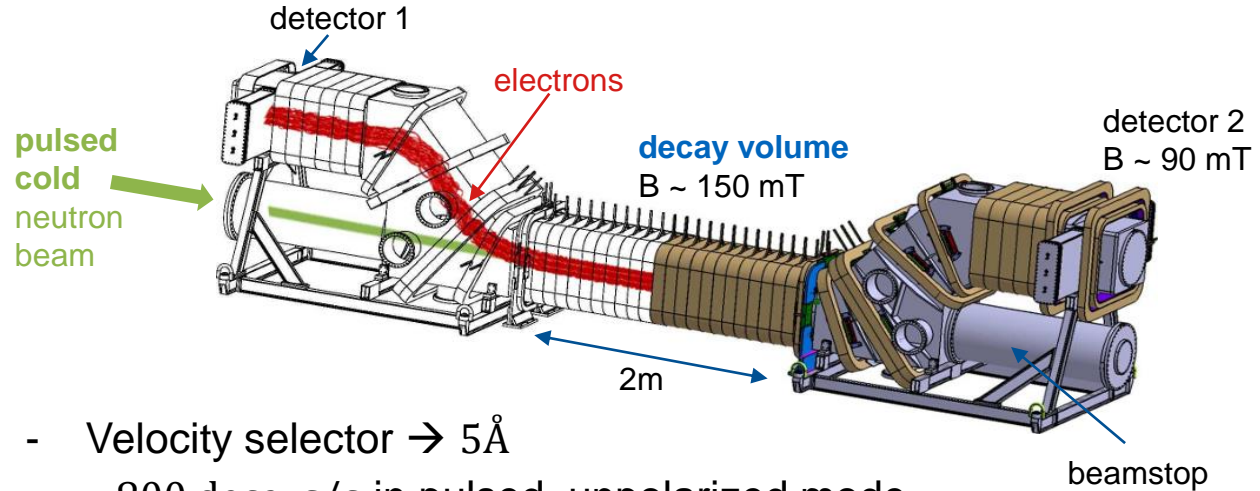
- Velocity selector $\rightarrow 5\text{\AA}$
- ~ 800 decays/s in pulsed, unpolarized mode

Months of preparation and setup } $\sim 10^9$ events
Three weeks of measurement }

PERKEO III



Originally built by University of Heidelberg, now operated by TUM, TU Vienna, HD & ILL.



- Velocity selector $\rightarrow 5\text{\AA}$
- ~ 800 decays/s in pulsed, unpolarized mode

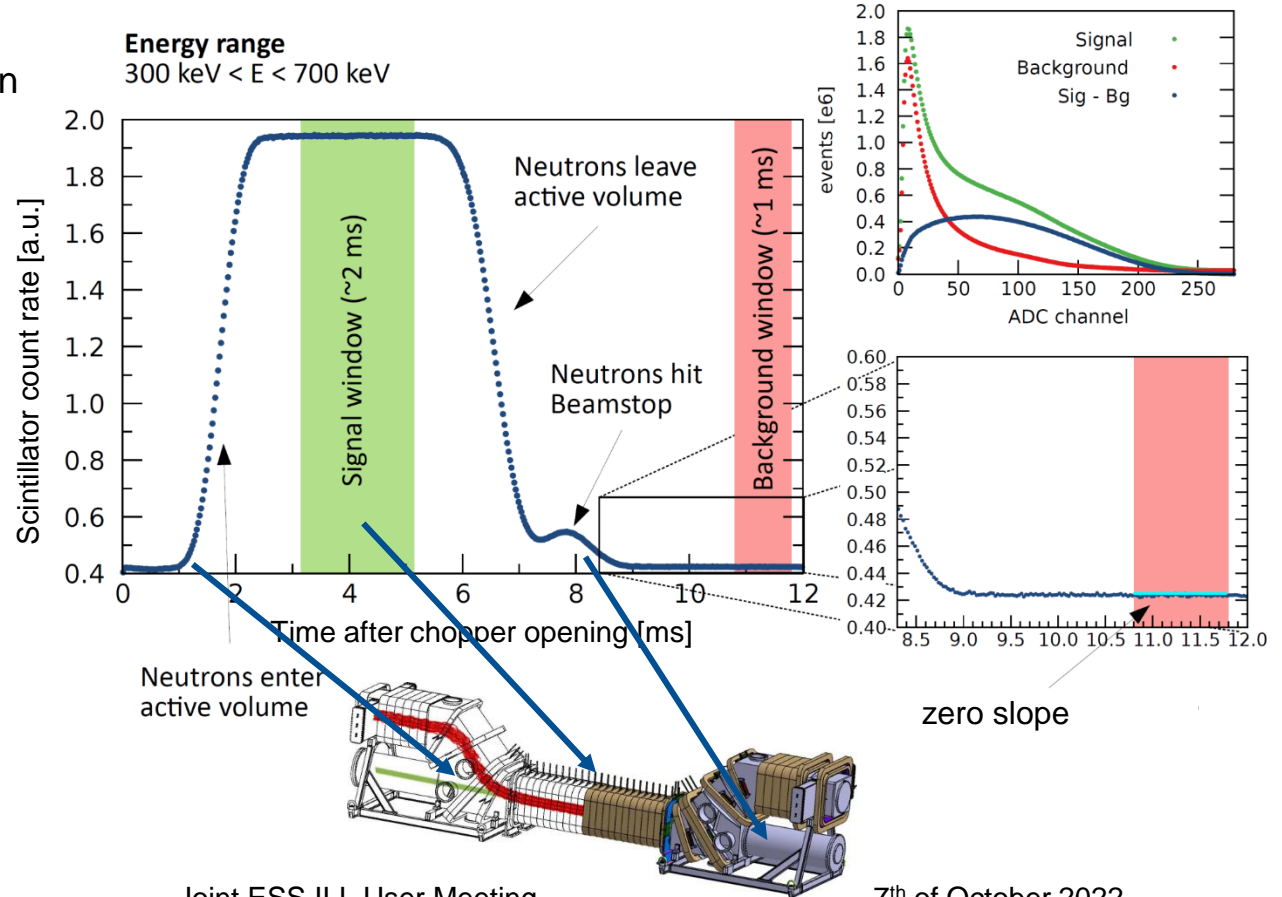
Months of preparation and setup } $\sim 10^9$ events
Three weeks of measurement }

\rightarrow Statistics ✓

Advantage of pulsed neutron beam

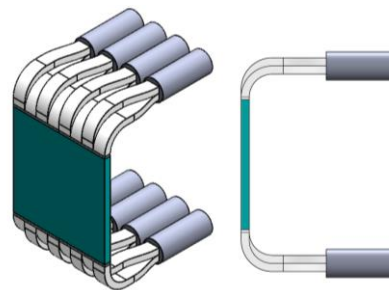
Same background condition in
 signal and background time
 window

→ Nearly perfect
 background
 subtraction

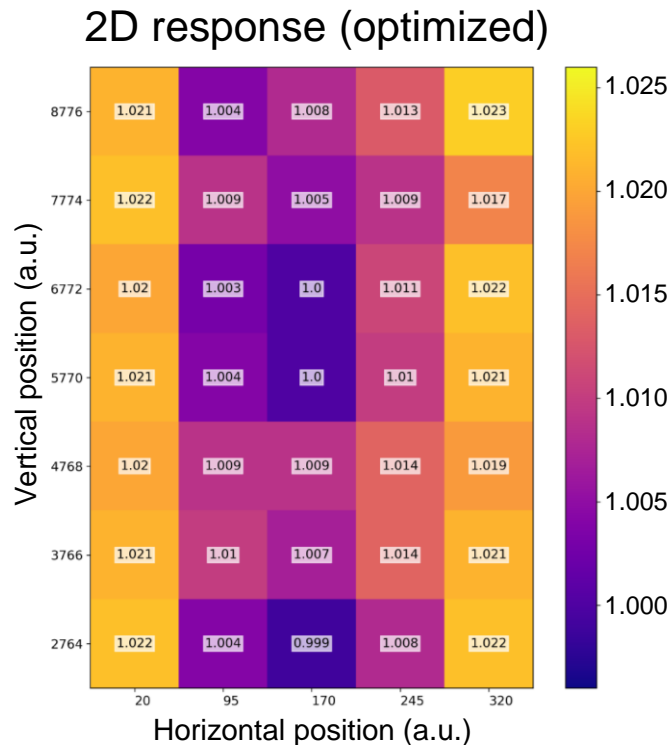


Improving systematics: Detector

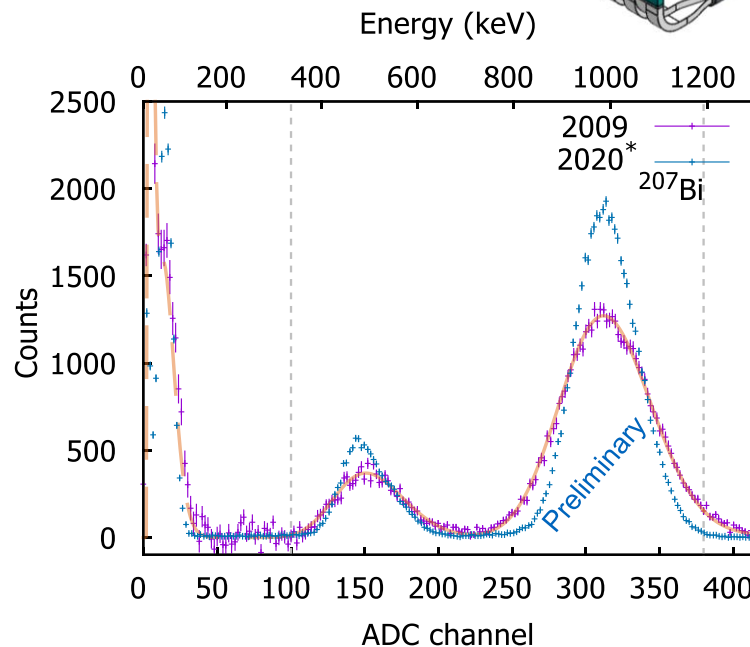
New detectors optimized for uniformity (~2%) and light output



Lightguides manufactured by Martin Aigner



Lamparth *et al.*, arXiv:2205.07625



**Factor of two
increase in
light output
gives higher
energy
resolution!**

Master Thesis K. Bernert 2019
Bachelor Thesis M. Anthony 2019

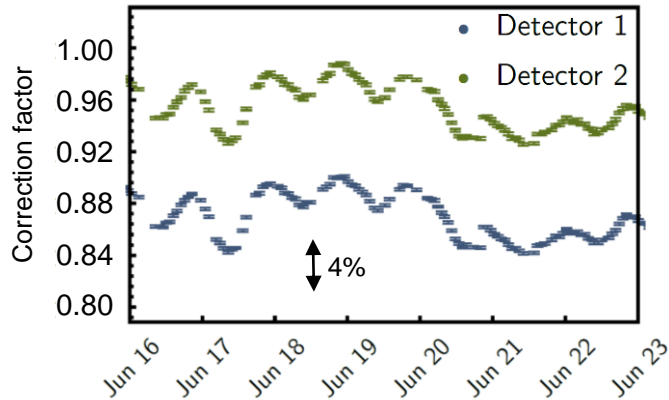
Improving systematics: Cooling

Shifts in temperature lead to shifts in PMT gain → Hourly drift measurements needed

Improving systematics: Cooling

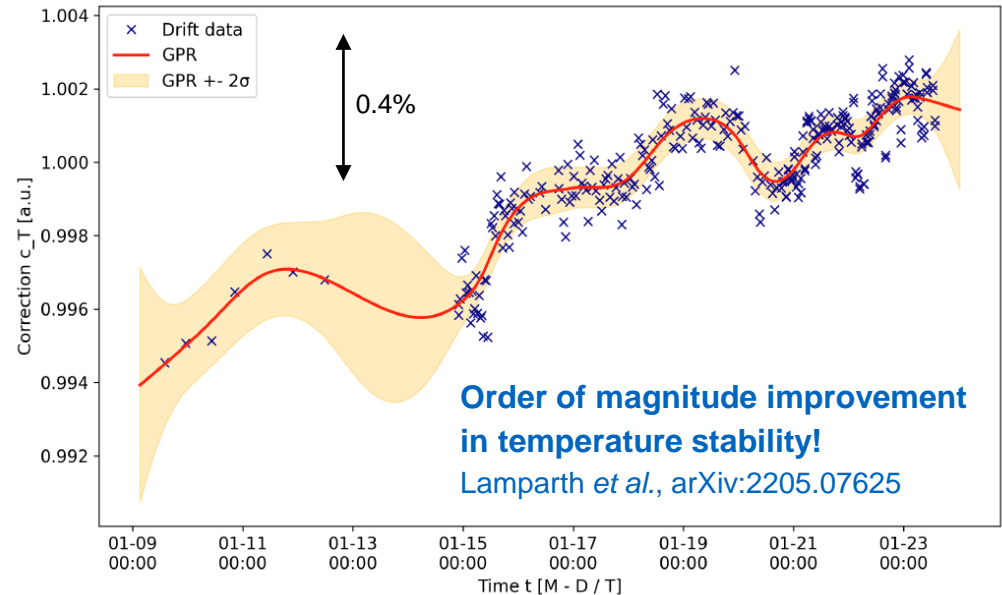
Shifts in temperature lead to shifts in PMT gain → Hourly drift measurements needed

Beta Asymmetry measurement in 2009



PhD Thesis, H. Saul (2018)

Fierz term measurement in 2019/20



Improving systematics: Cooling

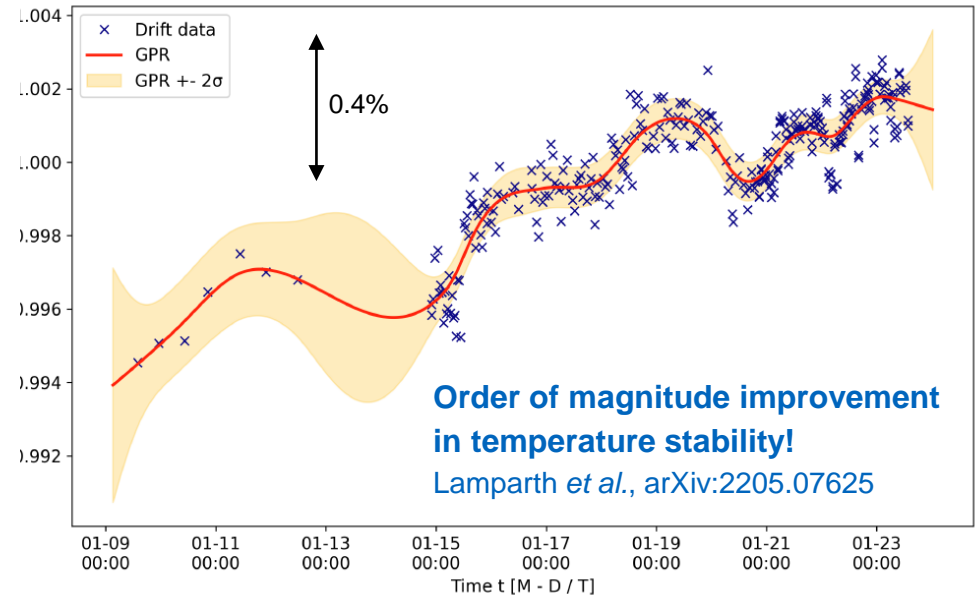
Shifts in temperature lead to shifts in PMT gain → Hourly drift measurements needed

New magnet cooling system helps to limit drift*

Fierz term measurement in 2019/20



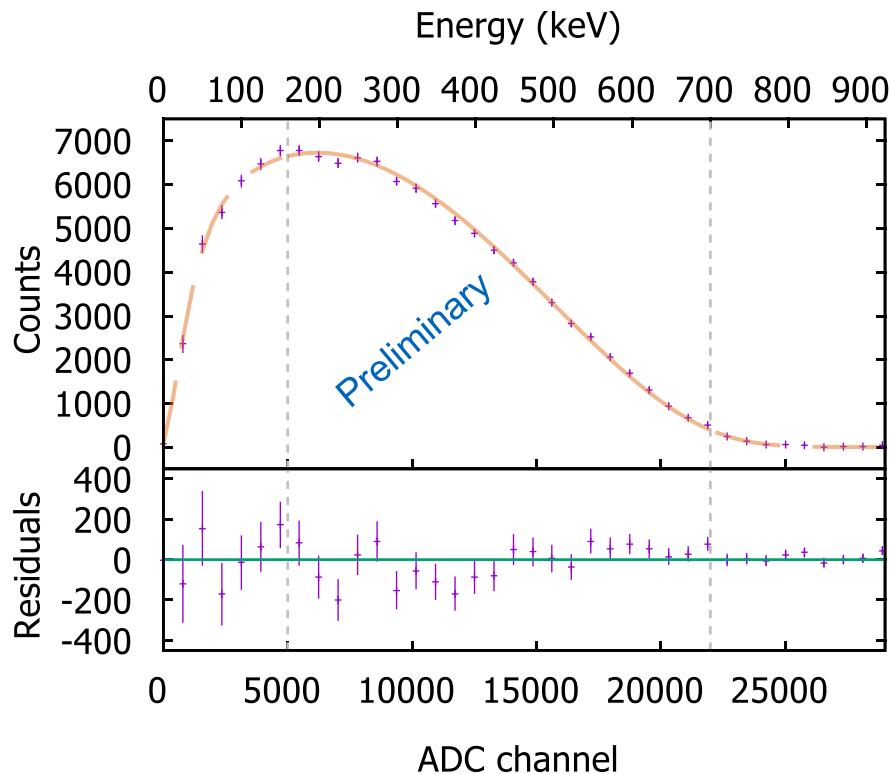
Measuring in the winter was helpful too ;)



* Dr. Jens Klenke, FRM II

A taste of the data...

5 min run



Outlook

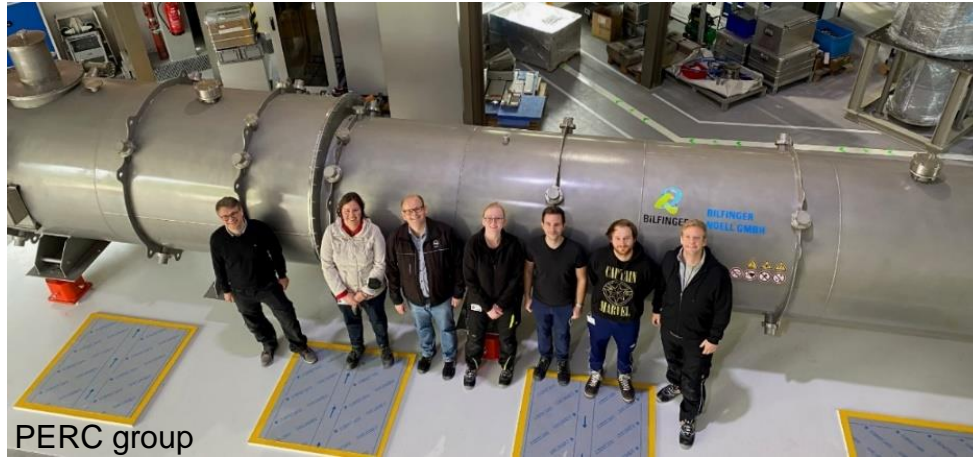
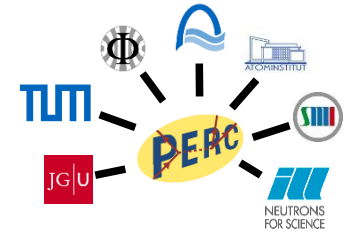


Analysis ongoing...

Outlook

Analysis ongoing...

Next generation experiment
is already under
construction at the FRM II



PERC (Proton Electron Radiation Channel)

- Longer decay volume \rightarrow higher event rate
- Magnetic filter for phase space selection
- Aims to improve measurements of A , C , a , b