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CEvNS light  $Z'$

Johan Rathsman

Light  $Z'$  models

Neutrino flux from  
 $\pi^+$  decay at rest

CEvNS

Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions

# Prospects for finding light $Z'$ from CEvNS at ESS

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FPNP@ESS, 2025-01-17



# Light $Z'$ models

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- Assume extra  $U(1)$  gauge symmetry giving additional  $Z'$  gauge boson
- B-L well known example

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$$\mathcal{L}_{\text{NC}}^{Z'} = - \sum_f \bar{\psi}_f (\gamma^\mu C_V^f + \gamma^\mu \gamma^5 C_A^f) \psi_f Z'_\mu$$

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Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions

In B-L models  $C_V^f = g_{Z'} Q'_f$  with  $Q'_u = Q'_d = 1/3$  and  $Q'_{\nu_\ell} = -1$ .

- here we assume fermion charges only constrained by anomaly conditions
- possible to explain so called Atomki anomaly with a 17 MeV  $Z'$  (X17)
- would modify normalisation and shape of nuclear recoil spectrum in CEvNS



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## Benchmark models allowed by current constraints

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Parameter	BM111	BM222	BM333	BM444	BM666
$g'$	$2.51 \times 10^{-5}$	$3.48 \times 10^{-5}$	$1.57 \times 10^{-5}$	$1.81 \times 10^{-5}$	$1.88 \times 10^{-5}$
$m_{Z'}$ [GeV]	0.0175	0.0174	0.0175	0.0166	0.0171
Neutrino flux from $\pi^+$ decay at rest					
$C_V^u$	$-2.49 \times 10^{-4}$	$-2.55 \times 10^{-4}$	$-2.48 \times 10^{-4}$	$-2.49 \times 10^{-4}$	$-2.54 \times 10^{-4}$
CEvNS					
$C_V^d$	$1.43 \times 10^{-4}$	$1.46 \times 10^{-4}$	$1.39 \times 10^{-4}$	$1.09 \times 10^{-4}$	$1.12 \times 10^{-4}$
Detector considerations					
$C_V^{\nu e}$	$1.38 \times 10^{-5}$	$1.39 \times 10^{-5}$	$1.10 \times 10^{-5}$	$-1.18 \times 10^{-5}$	$-1.13 \times 10^{-5}$
Statistics					
$C_V^{\nu \mu}$	$2.01 \times 10^{-5}$	$1.74 \times 10^{-5}$	$1.25 \times 10^{-5}$	$-1.27 \times 10^{-5}$	$-1.13 \times 10^{-5}$

Backgrounds and systematics

Summary and conclusions



# Neutrino flux from $\pi^+$ decay at rest

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Neutrino flux from  
 $\pi^+$  decay at rest

CEvNS

Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions

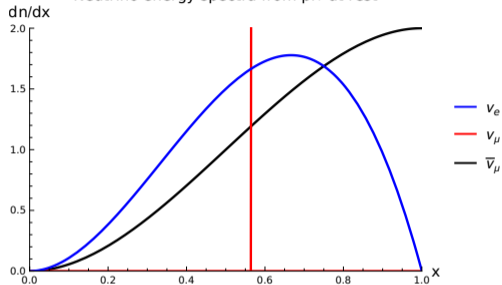
$$\frac{d\Phi_{\nu_e}}{dx} = \frac{rN_{\text{POT}}}{4\pi L^2} 12x^2(1-x)$$

$$\frac{d\Phi_{\bar{\nu}_\mu}}{dx} = \frac{rN_{\text{POT}}}{4\pi L^2} 2x^2(3-2x)$$

$$\frac{d\Phi_{\nu_\mu}}{dx} = \frac{rN_{\text{POT}}}{4\pi L^2} \delta(x-x_0)$$

- $x = \frac{2E_\nu}{m_\mu}$ ,  $0 < x < 1$
- $r$  nr of  $\pi^+$  per proton,
- $N_{\text{POT}}$  nr of protons on target,
- $L$  distance

Neutrino energy spectra from pi+ at rest



$$x_0 = \frac{m_{\pi^+}^2 - m_\mu^2}{m_\mu m_{\pi^+}} \approx 0.564.$$



# Nuclear recoil energy spectrum for CEvNS

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$$\frac{dN}{dy} = \frac{rN_{\text{POT}}}{4\pi L^2} \frac{G_F^2 m_\mu^2}{4\pi} \sum_{\nu_\ell} \left[ \frac{dn_{\nu_\ell}}{dy} \nu (Q_V^{\nu_\ell})^2 \right]$$

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Neutrino flux from  $\pi^+$  decay at rest

CEvNS

Detector considerations

Statistics

Backgrounds and systematics

Summary and conclusions

- $y = \frac{2ME_{\text{nr}}}{m_\mu^2}$ ,  $0 < y < 1$

- $M$  nucleus mass.

- $E_{\text{nr}}$  nuclear recoil energy

$$\frac{dn_{\nu_e}}{dy} \nu = \frac{1}{2} - 3y + 4y^{3/2} - \frac{3}{2}y^2$$

$$\frac{dn_{\nu_\mu}}{dy} \nu = \frac{1}{2} - 2y + 2y^{3/2} - \frac{1}{2}y^2 + \left( \frac{1}{2} - \frac{y}{2x_0^2} \right) \Theta \left( 1 - \frac{y}{x_0^2} \right)$$

$$Q_V^{\nu_\ell} = g_V^{p,\nu_\ell} Z F_{V,Z}(y) + g_V^{n,\nu_\ell} N F_{V,N}(y)$$



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## Klein Nystrand nuclear form factor

$F_{V,Z}(y)$  and  $F_{V,N}(y)$  nuclear form factors for protons and neutrons

$$F_{V,Z}(Q^2) = F_{V,N}(Q^2) = \frac{3}{(QR_A)^3} [\sin(QR_A) - QR_A \cos(QR_A)] \frac{1}{1 + Q^2}$$

$$Q^2 = 2ME_{\text{nr}} = ym_{\mu}^2, \quad R_A = 1.2A^{1/3} \text{ fm, and } a = 0.7 \text{ fm.}$$

Light Z' models

Neutrino flux from  
 $\pi^+$  decay at rest

CEvNS

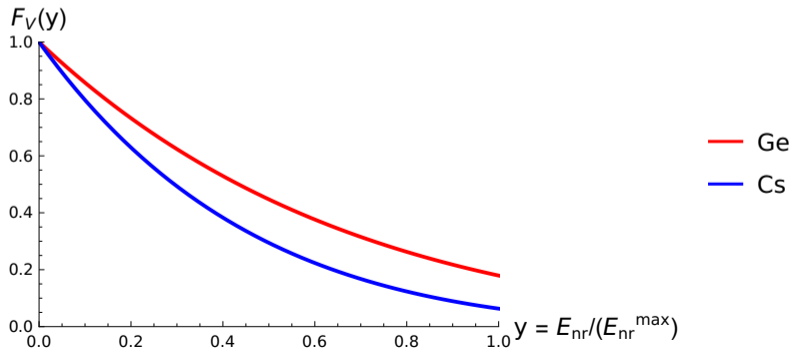
Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions

Form factor





# Vector couplings

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Vector couplings of neutrinos to protons and neutrons modified by  $Z'$  exchange

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$$g_V^{p,\nu_\ell} = g_{V,\text{SM}}^p + \frac{\sqrt{2}C_V^{\nu_\ell}(2C_V^u + C_V^d)}{G_F(y m_\mu^2 + m_{Z'}^2)}$$

Light  $Z'$  models

Neutrino flux from  
 $\pi^+$  decay at rest

$$g_V^{n,\nu_\ell} = g_{V,\text{SM}}^n + \frac{\sqrt{2}C_V^{\nu_\ell}(C_V^u + 2C_V^d)}{G_F(y m_\mu^2 + m_{Z'}^2)}$$

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Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions

for the SM we use the LO values

$$g_{V,\text{SM}}^p = (1 - 4 \sin^2 \theta_W)$$

$$g_{V,\text{SM}}^n = -1$$

with  $\sin^2 \theta_W = 0.24$



# Z' signal ratio compared to SM

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Using  $F_{V,Z}(y) = F_{V,N}(y)$  the signal ratio compared to the SM becomes

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$$\left[ \sum_{\nu_\ell} \frac{dn_{\nu_\ell}}{dy} \left( -1 + \frac{\sqrt{2}C_V^{\nu_\ell} [N(C_V^u + 2C_V^d) + Z(2C_V^u + C_V^d)]}{[N - (1 - 4 \sin^2 \theta_W)Z] G_F m_\mu^2 (y + m_{Z'}^2/m_\mu^2)} \right)^2 \right] / \left[ \sum_{\nu_\ell} \frac{dn_{\nu_\ell}}{dy} \right]$$

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Neutrino flux from  $\pi^+$  decay at rest

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for  $B - L$  model this reduces to

Detector

considerations

Statistics

Backgrounds and systematics

Summary and conclusions

$$\left( 1 + \frac{\sqrt{2}g_{Z'}^2 (N + Z)}{[N - (1 - 4 \sin^2 \theta_W)Z] G_F m_\mu^2 (y + m_{Z'}^2/m_\mu^2)} \right)^2$$





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Current limits on B-L model from Atzori Corona et al (2202.11002):

$$g_{Z'} < 5 \times 10^{-5} \text{ for } m_{Z'} = 0.017 \text{ GeV}$$

- will use this as a reference model/parameter value

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Light  $Z'$  models

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 $\pi^+$  decay at rest

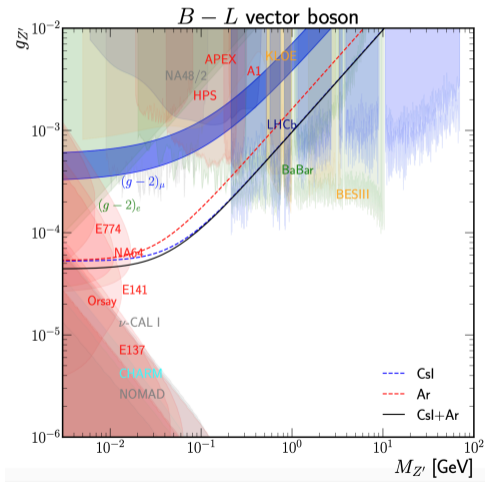
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Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions





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The propagator  $\frac{1}{y + m_{Z'}^2/m_\mu^2}$  affects the shape of the spectrum

Ex: B-L model with  $g_{Z'} = 5 \times 10^{-5}$  for  $m_{Z'} = 4.25, 8.5, 17, 34$  MeV

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 $\pi^+$  decay at rest

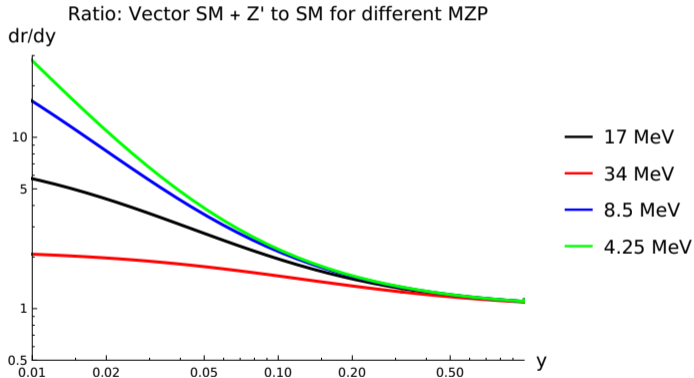
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Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions



note: independent of nuclear mass



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# Signal ratio when integrating over $0.01 < y < 0.31$

$$\text{Effective couplings: } C_{\text{eff}}^{\nu\ell} = \frac{C_V^{\nu\ell} [N(C_V^u + 2C_V^d) + Z(2C_V^u + C_V^d)]}{[N - (1 - 4\sin^2\theta_W)Z]}$$

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Neutrino flux from  
 $\pi^+$  decay at rest

CEvNS

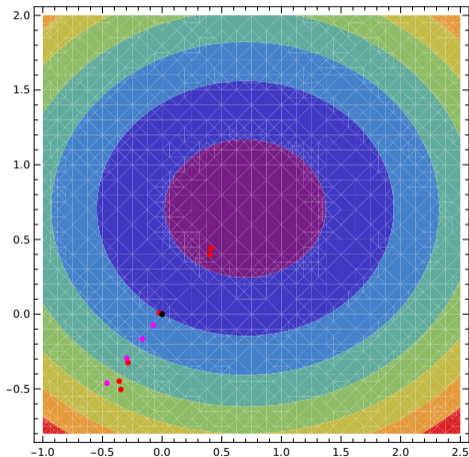
Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions

$C_{\text{eff}}^{\nu\mu} [10^{-8}]$



$C_{\text{eff}}^{\nu e} [10^{-8}]$

signal ratio



black: SM

magenta: B-L

$g_{Z'} =$   
 $2, 3, 4, 5 \times 10^{-5}$

red: X17 BMs



# Detector considerations

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Germanium detector as an example (inspired by Dreden-II exp)

Quenching factor: conversion of nuclear recoil to ionisation signal

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$$y_{\text{ion}} = Q(y)y$$

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Neutrino flux from  
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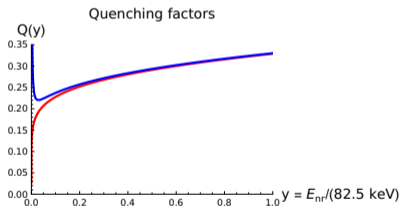
CEvNS

Detector  
considerations

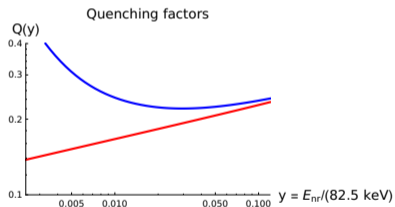
Statistics

Backgrounds and  
systematics

Summary and  
conclusions



— Lindhard  $k=0.157$   
— Lindhard + Mod



Modifications added to span uncertainties at small recoil energies

Threshold  $\sim 800 \text{ eV} \Rightarrow y \gtrsim 0.01$



# Detector considerations, cont'd

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## The detector resolution

$$R(y_{\text{rec}}, y_{\text{ion}}) = \frac{1}{\sigma\sqrt{2\pi}} \exp\left[-\frac{(y_{\text{rec}} - y_{\text{ion}})^2}{2\sigma^2}\right]$$

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Neutrino flux from  $\pi^+$  decay at rest

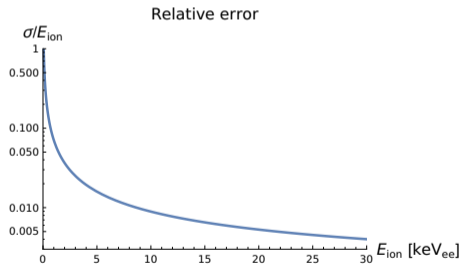
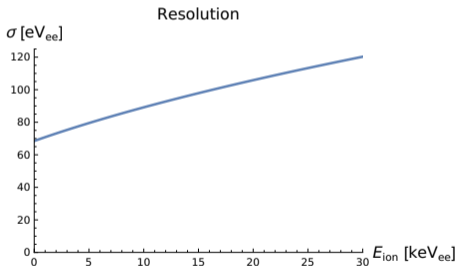
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**Detector considerations**

Statistics

Backgrounds and systematics

Summary and conclusions



$$\text{Threshold: } y_{\text{rec}} > \frac{0.200 \text{ keV}}{82.5 \text{ keV}} = 0.0024$$



# Statistics

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Flux factor  $\frac{rN_{\text{POT}}}{4\pi L^2}$  assuming 5000 hours per year of beam on target

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Detector considerations

Statistics

Backgrounds and systematics

Summary and conclusions

Parameter	Far, $L=25$ m		
	Low	Med	High
Beam energy [GeV]	0.84	0.84	2.0
Yield $r$	0.08	0.08	0.3
Power [MW]	0.80	2.0	5.0
$N_{\text{POT}}$ [year $^{-1}$ ]	$1.1 \times 10^{23}$	$2.7 \times 10^{23}$	$2.7 \times 10^{23}$
Time in years	1	3	3
Flux factor [GeV $^2$ ]	$4.3 \times 10^{-14}$	$3.2 \times 10^{-13}$	$1.2 \times 10^{-12}$

Also assuming  $m_{\text{target}} = 20$  kg





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Neutrino flux from  
 $\pi^+$  decay at rest

CEvNS

Detector  
considerations

Statistics

**Backgrounds and  
systematics**

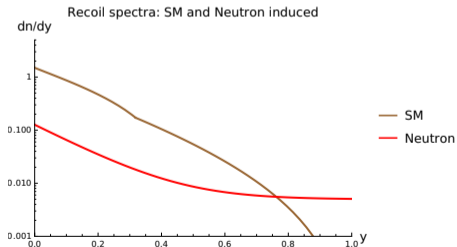
Summary and  
conclusions

# Backgrounds and systematics

## Backgrounds

- Reducible: neutron induced nuclear recoils
  - prompt
  - ambient - here the duty factor helps
- Irreducible
  - Standard Model

Neutron background at  $L = 25$  m with 50 cm polyethylene neutron moderator



Based on C. M. Lewis (PhD thesis U Chicago) (rescaled CsI to Ge)





# Backgrounds and systematics, cont'd

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## Systematic uncertainties

- neutrino flux - uncertainties in  $r$  - assume 10 %
- quenching factor

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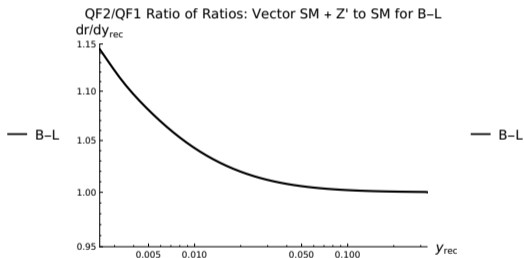
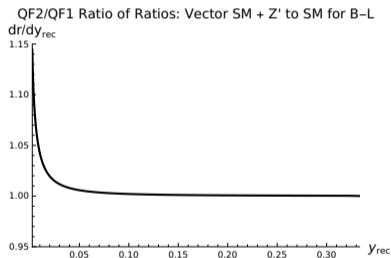
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Detector considerations

Statistics

Backgrounds and systematics

Summary and conclusions



- small for  $y_{rec} > 0.02$
- up to 15 % at threshold  $y_{rec} = 0.0024$  (0.200 keV<sub>ee</sub>)

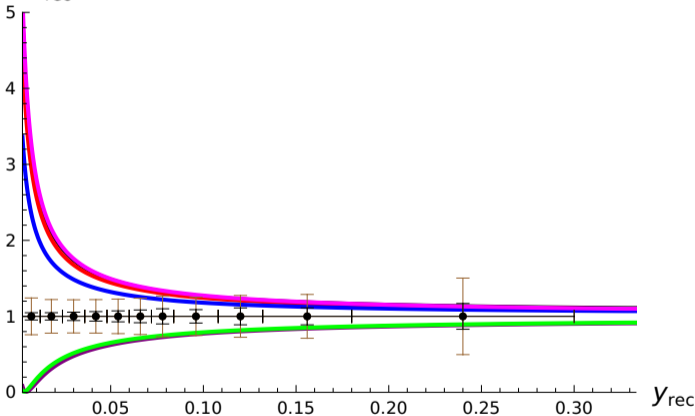


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# Resulting ratio to SM

Smearing QF1 Ratio: Vector SM + Z' to SM for different BM

$dR/dy_{rec}$



- BM111
- BM222
- BM333
- BM444
- BM666
- B-L
- SM

statistical error (in black)  $1/\sqrt{N_{SM}} \sim 5\%$

systematic errors:  $R_{QF} + R_{flux} + R_{neutron}$

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CEvNS

Detector considerations

Statistics

Backgrounds and systematics

Summary and conclusions



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Light  $Z'$  models

Neutrino flux from  
 $\pi^+$  decay at rest

CEvNS

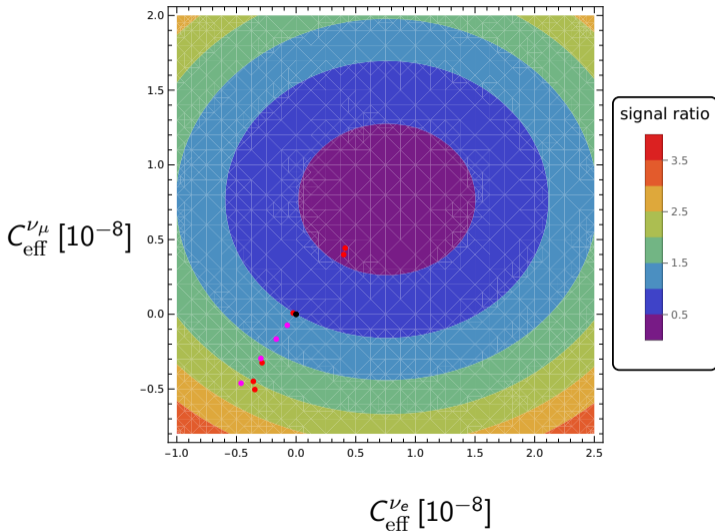
Detector  
considerations

Statistics

Backgrounds and  
systematics

Summary and  
conclusions

Integrating over  $0.0024 < y_{\text{rec}} < 0.1$  gives overall signal strength



black: SM

magenta: B-L

$g_{Z'} =$   
 $2, 3, 4, 5 \times 10^{-5}$

red: X17 BMs



# Summary and conclusions

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Light  $Z'$  models

Neutrino flux from  
 $\pi^+$  decay at rest

CEvNS

Detector  
considerations

Statistics

Backgrounds and  
systematics

**Summary and  
conclusions**

- CEvNS very promising for finding light  $Z'$  gauge boson
- For  $m_{Z'} \sim 20$  MeV the shape could also be used
- Current limits allow for a factor 2 up or down compared to SM cross-section
- Currently systematics limiting factor
- few percent statistical uncertainty possible to reach with 5 MW